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Superoscillating Sequences and Supershifts for Families of Generalized Functions

F. Colombo¹ · I. Sabadini¹ · D. C. Struppa² · A. Yger³

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Abstract

We construct a large class of superoscillating sequences, more generally of \mathcal{F} -supershifts, where \mathcal{F} is a family of smooth functions in (t, x) (resp. distributions in (t, x) , or hyperfunctions in x depending on the parameter t) indexed by $\lambda \in \mathbb{R}$. The frame in which we introduce such families is that of the evolution through Schrödinger equation $(i\partial/\partial t - \mathcal{H}(x))(\psi) = 0$ ($\mathcal{H}(x) = -(\partial^2/\partial x^2)/2 + V(x)$), V being a suitable potential). If $\mathcal{F} = \{(t, x) \mapsto \varphi_\lambda(t, x); \lambda \in \mathbb{R}\}$, where φ_λ is evolved from the initial datum $x \mapsto e^{i\lambda x}$, \mathcal{F} -supershifts will be of the form $\{\sum_{j=0}^N C_j(N, a)\varphi_{1-2j/N}\}_{N \geq 1}$ for $a \in \mathbb{R} \setminus [-1, 1]$, taking $C_j(N, a) = \binom{N}{j}(1+a)^{N-j}(1-a)^j/2^N$. Our results rely on the fact that *integral operators of the Fresnel type* govern, as in optical diffraction, the evolution through the Schrödinger equation, such operators acting continuously on the weighted algebra of entire functions $\text{Exp}(\mathbb{C})$. Analyzing in particular the quantum harmonic oscillator case forces us, in order to take into account singularities of the evolved datum that occur when the stationary phasis in the Fresnel operator vanishes, to enlarge the notion of \mathcal{F} -supershift, \mathcal{F} being a family of C^∞ functions or distributions in (t, x) , to that where \mathcal{F} is a family of hyperfunctions in x , depending on t as a parameter.

Keywords Superoscillations · Infinite order differential operators · Schrödinger equation · Generalized functions

Mathematics Subject Classification 42A16 · 30D15 · 46F15

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1 Introduction

The Aharonov–Berry superoscillations are band-limited functions that can oscillate faster than their fastest Fourier component. These functions (or sequences, to be more precise) appear in many settings, both in classical phenomena, as well as in the study of weak measurements, [7, 14, 15, 29]. A thorough review of the various situations in which superoscillations arise is [30]. Not surprisingly, the literature related to superoscillations is very large, and without claiming completeness, we mention [23–25, 27, 28, 36, 37, 43, 45]. Quite recently, this class of functions has been investigated from the mathematical point of view, see [1, 2, 8–12, 16, 18–20, 31, 33, 34] and the monograph [13]. For a unified approach to Schrödinger evolution of superoscillations and supershifts see the recent paper [3]. Finally, we mention that different field equations has been recently considered such as Klein–Gordon see [5], and that superoscillations also appear in Schur analysis, see [17].

In this paper we will extend naturally some of the ideas developed in those articles, and we will discover exciting new features that need to be considered when we apply the theory of superoscillations to the study of specific Schrödinger equations, in which singularities naturally arise.

In order to set up the stage, let $a > 1$ be a real number. The archetypical superoscillatory sequence is the sequence of complex valued functions $\{F_N(x, a)\}_{N \geq 1}$ defined on \mathbb{R} by

$$F_N(x, a) = \left(\cos\left(\frac{x}{N}\right) + ia \sin\left(\frac{x}{N}\right) \right)^N. \quad (1.1)$$

It is immediate to see that if we fix $x \in \mathbb{R}$, and we let N go to infinity, we obtain that

$$\lim_{N \rightarrow \infty} F_N(x, a) = e^{iax}.$$

However, by expanding the trigonometric functions $\sin x$ and $\cos x$ through Euler's formula, it is also easy to see that in fact we can rewrite $F_N(x, a)$ as

$$F_N(x, a) = \sum_{j=0}^N C_j(N, a) e^{i(1-2j/N)x},$$

where

$$C_j(N, a) = \binom{N}{j} \left(\frac{1+a}{2}\right)^{N-j} \left(\frac{1-a}{2}\right)^j, \quad (1.2)$$

and $\binom{N}{j}$ denotes the binomial coefficient. This simple observation is significant because we now have a sequence of functions whose frequencies $(1 - 2j/N)$ are bounded by one, and yet their limit has the arbitrarily large frequency a . This explains why the sequence $\{F_N(x, a)\}_{N \geq 1}$ is called *superoscillatory*.

There is also a different way to conceptualize this phenomenon by saying that the map that associates to a real number $a \in \mathbb{R} \setminus [-1, 1]$ the sequence $\{F_N(x, a)\}$ is a supershift with respect to the map which associates to the real parameter λ the complex function $e^{i\lambda x}$. Though this is simply a different way to restate what we said before, it indicates a change of perspective since the supershift is now associated to maps whose variables are real numbers and whose range are spaces of functions. This approach will be made precise in Definition 5.1.

Most of the work done in recent years has been devoted to two fundamental questions, inspired by the physical considerations that led to the discovery of superoscillatory phenomena. The first question asks how to construct larger families of superoscillating sequences, and the second question asks whether superoscillations persist once they are propagated according to the Schrödinger equation with a given potential. As it turns out, the two questions are intimately connected, and this paper explores further this connection. To begin with, it is easy to show that if $P(z)$ is a polynomial of one complex variable, then the sequence

$$\psi_{P,N}(t, x, \lambda) = \sum_{j=0}^N C_j(N, \lambda) e^{i(1-2j/N)x} e^{itP(1-2j/N)},$$

with t and x real variables, is a solution of the Cauchy problem for the modified Schrödinger equation

$$\left[i \frac{\partial}{\partial t} + P \left(-i \frac{\partial}{\partial x} \right) \right] (\psi_{P,N}(t, x, \lambda)) = 0,$$

with the initial condition

$$\psi_{P,N}(0, x, \lambda) = F_N(x, \lambda).$$

In Sect. 3, however, we will show that this new sequence $\{\psi_{P,N}(t, x, \lambda)\}$ converges, with all of its derivatives in both t and x , to the suitable derivative of $e^{i(\lambda x + tP(\lambda))}$. We are therefore able to show that this process allows us to construct a large class of superoscillating sequences of the form $\{x \mapsto \psi_{P,N}(t, x, a)\}_{N \geq 1}$ from the original “superoscillating” sequence $\{x \mapsto F_N(x, a)\}_{N \geq 1}$. In these sequences the superoscillating convergence property propagates through any differential operator in t, x , the convergence being uniform on any compact subset of the real plane. Note that this process as well allows us to extend the notion of supershift to functions that associate to a real parameter λ a function of two variables t and x .

The more significant situation, however, arises when one considers a general Schrödinger operator of the form

$$i \frac{\partial}{\partial t} + \frac{1}{2} \frac{\partial^2}{\partial x^2} + V(x)$$

with a suitable real potential V and associated Green function $G_V(t, x, 0, x')$ such that the response, at time t , to the initial value $e^{i\lambda x}$ for $t = 0$, is given by

$$\varphi_\lambda(t, x) = \int_{\mathbb{R}} G_V(t, x, 0, x') e^{i\lambda x'} dx'$$

for $\lambda \in \mathbb{R}$ and (t, x) in the phase domain. Note that we need to interpret the integral as a regularized integral on \mathbb{R} in a sense that will be clarified in Sect. 4). In this case we will find sufficient conditions that ensure that the integral operator

$$T(f(x'))(t, x) = \int_{\mathbb{R}} G_V(t, x, 0, x') f(x') dx'$$

is such that for any $\lambda \in \mathbb{R}$,

$$T\left(\sum_{j=0}^N C_j(N, \lambda) e^{i(1-2j/N)x'}\right)(t, x) = \sum_{j=0}^N C_j(N, \lambda) \varphi_{1-2j/N}(t, x)$$

converges to

$$T(e^{i\lambda x'})(t, x) = \varphi_\lambda(t, x)$$

locally uniformly in some open subset \mathcal{U} of the phase space (in $\mathbb{R}_{t,x}^2$) on which V is smooth and which is entirely determined by the explicit expression of the Green function G_V (Theorem 5.1). In such a situation the sequence defined on \mathcal{U} by

$$\left\{ \sum_{j=0}^N C_j(N, a) \varphi_{1-2j/N} \right\}$$

(where $a \in \mathbb{R} \setminus [-1, 1]$) is a supershift for $(\varphi_\lambda)|_{\mathcal{U}}$. Moreover, for any $\lambda \in \mathbb{R}$, $\varphi_\lambda \in \mathcal{C}^\infty(\mathcal{U}, \mathbb{C})$ and for any $(\mu, \nu) \in \mathbb{N}^2$,

$$\frac{\partial^{\mu+\nu}}{\partial t^\mu \partial x^\nu} \left(\sum_{j=0}^N C_j(N, \lambda) \varphi_{1-2j/N} \right) \xrightarrow{N \rightarrow \infty} \frac{\partial^{\mu+\nu}}{\partial t^\mu \partial x^\nu} (\varphi_\lambda)$$

locally uniformly in \mathcal{U} for any $\lambda \in \mathbb{R}$, in particular $\lambda = a \in \mathbb{R} \setminus [-1, 1]$.

Interesting situations occur when $\lambda \in \mathbb{R} \mapsto \varphi_\lambda$ makes sense as a continuous map from \mathbb{R} into $\mathcal{D}'(\mathcal{U}', \mathbb{C})$ for some open subset $\mathcal{U}' \supseteq \mathcal{U}$ in the phase space. Such is the case in the example of the quantum harmonic oscillator, where $V(x) = x^2/2$, the phase space is $\mathbb{R}^{+*} \times \mathbb{R}$ and

$$\mathcal{U} = (\mathbb{R}^+ \times \mathbb{R}) \setminus \{(k\pi/2, x); k \in \mathbb{N}, x \in \mathbb{R}\} \subset \mathcal{U}' = \mathbb{R}^{+*} \times \mathbb{R}.$$

In such case, given $k' \in \mathbb{N}$ and $x_0 \in \mathbb{R}$, it is impossible to interpret

$$\left\{ \left(\sum_{j=0}^N C_j(N, a) \varphi_{1-2j/N} \right)_{\text{about } ((2k'+1)\pi/2, x_0)} \right\} \tag{1.3}$$

(when $a \in \mathbb{R} \setminus [-1, 1]$) as a supershift for $\lambda \mapsto (\varphi_\lambda)_{\text{about } ((2k'+1)\pi/2, x_0)}$ (all maps being considered here as distribution-valued about $((2k' + 1)\pi/2, x_0)$), while it is possible to do so about a point $(k''\pi, x_0)$, where $k'' \in \mathbb{N}^*$. In order to interpret (1.3) as a supershift for $\lambda \mapsto (\varphi_\lambda)_{\text{about } ((2k'+1)\pi/2, x_0)}$, one needs to consider $(\varphi_\lambda)_{\text{about } ((2k'+1)\pi/2, x_0)}$ as a hyperfunction (in t) times a distribution (in x) instead of distribution in (t, x) . We will discuss such questions in Sect. 6.

The plan of the paper is the following: the paper contains five sections, besides this introduction. In Sect. 2 we introduce the spaces $A_p(\mathbb{C})$, $A_{p,0}(\mathbb{C})$, and we define some infinite order differential operators with nonconstant coefficients which will play a crucial role to prove our main results. In Sect. 3 we recall the definition of generalized Fourier sequence and (complex) superoscillating sequence in one and several variables together with some examples; we then study two Cauchy–Kowalevski problems (one of which of Schrödinger type) and we show that superoscillations persist in time. In Sect. 4 we address the problem of explaining the process of regularization of formal Fresnel-type integrals which is a necessary step to obtain further results in the paper. Fresnel-type integrals are shown to be continuous on $A_1(\mathbb{C})$ in Sect. 5, in which we also treat a Cauchy problem for the Schrödinger equation with centrifugal potential and also for the quantum harmonic oscillator. Finally, in Sect. 6, we investigate the evolution of superoscillating initial data with respect to the notion of supershift for the quantum harmonic oscillator, and we focus on singularities. It is interesting to note that in this case one needs to extend the concept of supershift in the case of hyperfunctions.

Notations. We use the notations with capital letters $Z, d/dZ, W, d/dW, \check{Z}$ in the expressions of formal differential operators, besides the usual notation z for the complex variable and t (time) x, x' (space) real variables.

2 On Continuity of Some Convolution Operators

Let f be a non-constant entire function of a complex variable z . We define

$$M_f(r) = \max_{|z|=r} |f(z)|, \quad \text{for } r \geq 0.$$

The non-negative real number ρ defined by

$$\rho = \limsup_{r \rightarrow \infty} \frac{\ln \ln M_f(r)}{\ln r}$$

is called the *order* of f . If ρ is finite then f is said to be of *finite order* and if $\rho = \infty$ the function f is said to be of *infinite order*.

In the case f is of finite order we introduce the non-negative real number

$$\sigma = \limsup_{r \rightarrow \infty} \frac{\ln M_f(r)}{r^\rho},$$

and call it the *type* of f . If $\sigma \in (0, \infty)$ we say f is of *normal type*, while we say it is of *minimal type* if $\sigma = 0$ and of *maximal type* if $\sigma = \infty$. Constant entire functions are considered of minimal type and order zero. In the sequel we will extensively make use of *weighted spaces* $A_p(\mathbb{C})$ or $A_{p,0}(\mathbb{C})$ of entire functions whose definition follows ; such spaces are classical, see e.g. [22, 47].

Definition 2.1 Let p be a strictly positive number. We define the space $A_p(\mathbb{C})$ as the \mathbb{C} -algebra of entire functions such that there exists $B > 0$ such that

$$\sup_{z \in \mathbb{C}} (|f(z)| \exp(-B|z|^p)) < +\infty.$$

The space $A_{p,0}(\mathbb{C})$ consists of those entire functions such that

$$\forall \varepsilon > 0, \sup_{z \in \mathbb{C}} (|f(z)| \exp(-\varepsilon|z|^p)) < +\infty.$$

To define a topology in these spaces we follow [22, Section 2.1]. For $p > 0$, $B > 0$ and for any entire function f , we set

$$\|f\|_B := \sup_{z \in \mathbb{C}} \{|f(z)| \exp(-B|z|^p)\}.$$

Let $A_p^B(\mathbb{C})$ denote the \mathbb{C} -vector space of entire functions satisfying $\|f\|_B < \infty$. Then $\|\cdot\|_B$ defines a norm on $A_p^B(\mathbb{C})$ so that $(A_p^B(\mathbb{C}), \|\cdot\|_B)$ is a Banach space and the natural inclusion mapping $A_p^{B'} \hookrightarrow A_p^B$ (when $0 < B \leq B'$) is a compact operator from $(A_p^{B'}(\mathbb{C}), \|\cdot\|_{B'})$ into $(A_p^B(\mathbb{C}), \|\cdot\|_B)$. For any sequence $\{B_n\}_{n \geq 1}$ of positive numbers, strictly increasing to infinity, we can introduce an LF-topology on $A_p(\mathbb{C})$ given by the inductive limit

$$A_p(\mathbb{C}) := \varinjlim A_p^{B_n}(\mathbb{C}).$$

Since this topology is stronger than the topology of the pointwise convergence, it is independent of the choice of the sequence $\{B_n\}_{n \geq 1}$. Thus, in this inductive limit topology, given f and a sequence $\{f_N\}_{N \geq 1}$ in $A_p(\mathbb{C})$, we say that $f_N \rightarrow f$ in $A_p(\mathbb{C})$ if and only if there exists $n \in \mathbb{N}^*$ such that $f, f_N \in A_p^{B_n}(\mathbb{C})$ for all $N \in \mathbb{N}^*$, and $\|f_N - f\|_{B_n} \rightarrow 0$ for $N \rightarrow \infty$. The topology on $A_{p,0}(\mathbb{C})$ is given as the projective limit

$$A_{p,0}(\mathbb{C}) := \varprojlim A_p^{\varepsilon_n}(\mathbb{C})$$

where $\{\varepsilon_n\}_{n \geq 1}$ is a strictly decreasing sequence of positive numbers converging to 0. It can be proved, see [22, Section 6.1], that $A_p(\mathbb{C})$ and $A_{p,0}(\mathbb{C})$ are respectively a DFS space and an FS space. When $p > 1$, $A_{p,0}(\mathbb{C})$ is the strong dual of $A_{p'}(\mathbb{C})$ (where $1/p + 1/p' = 1$), the duality being realized as

$$\mu \in (A_{p'}(\mathbb{C}))' \mapsto \left[\text{Fourier - Borel Transform of } \mu : w \in \mathbb{C} \mapsto \mu_z(e^{-zw}) \right] \in A_{p,0}(\mathbb{C}).$$

In the extreme case $p = 1$, $A_1(\mathbb{C})$ (also denoted as $\text{Exp}(\mathbb{C})$) is isomorphic to the space $\widehat{H(\mathbb{C})}$ of analytic functionals, the duality being realized as

$$T \in \widehat{H(\mathbb{C})} \mapsto \left[\text{Fourier - Borel Transform of } T : w \in \mathbb{C} \mapsto T_z(e^{-zw}) \right] \in A_1(\mathbb{C}).$$

Here $H(\mathbb{C})$ is equipped with its usual topology of uniform convergence on any compact subset.

The following result is an immediate consequence of the definition of the topology in the spaces $A_p(\mathbb{C})$ for $p > 0$.

Proposition 2.1 *Let $f = \{f_N\}_{N \geq 1}$ be a sequence of elements in $A_p(\mathbb{C})$. The two following assertions are equivalent:*

- the sequence f converges towards 0 in $A_p(\mathbb{C})$;
- the sequence f converges towards 0 in $H(\mathbb{C})$ and there exists $A_f \geq 0$ and $B_f \geq 0$ such that

$$\forall N \in \mathbb{N}^*, \quad \forall z \in \mathbb{C}, \quad |f_N(z)| \leq A_f e^{B_f |z|^p}. \tag{2.1}$$

Proof The first assertion means that there exists $B > 0$ with $\lim_{N \rightarrow \infty} \|f_N\|_B = 0$ (in particular $\|f_N\|_B \leq 1$ for $N \geq N_1$), which implies that the sequence f converges to 0 in $H(\mathbb{C})$ and that $|f_N(z)| \leq A e^{B|z|^p}$ with B and $A = \sup(\tilde{A}_1, \dots, \tilde{A}_{N_1}, 1)$ independent of N ($\tilde{A}_j = \sup_{\mathbb{C}}(|f_j(z)| e^{-B|z|^p})$ for $j = 1, \dots, N_1$). Conversely, assume that the second assertion holds and take $B > B_f$, so that, given $\varepsilon > 0$, there exists $R_\varepsilon > 0$ such that

$$\forall N \in \mathbb{N}^*, \quad \sup_{|z| \geq R_\varepsilon} |f_N(z)| e^{-B|z|^p} \leq A_f e^{(B_f - B)R_\varepsilon^p} < \varepsilon.$$

On the other hand, since f converges to 0 uniformly on any compact subset of \mathbb{C} , in particular on $\overline{D(0, R_\varepsilon)}$, there exists $N_\varepsilon \in \mathbb{N}^*$ such that

$$N \geq N_\varepsilon \implies \sup_{|z| \leq R_\varepsilon} |f_N(z)| e^{-B|z|^p} \leq \sup_{|z| \leq R_\varepsilon} |f_N(z)| < \varepsilon.$$

Therefore $\sup_{N \geq N_\varepsilon} \|f_N\|_B < \varepsilon$ and the sequence f converges to 0 in $A_p(\mathbb{C})$. □

To prove our main results we need an important lemma that characterizes entire functions in $A_p(\mathbb{C})$ in terms of the behaviour of their Taylor development, see Lemma 2.2 in [18].

Lemma 2.1 *The entire function $f : z \mapsto \sum_{j=0}^{\infty} f_j z^j$ belongs to $A_p(\mathbb{C})$ if and only if there exists $C = C_f > 0$ and $b = b_f > 0$ such that $f \in A_p^{C,b}(\mathbb{C})$, where*

$$A_p^{C,b}(\mathbb{C}) = \left\{ \sum_{j=0}^{\infty} f_j z^j \in A_p(\mathbb{C}); \forall j \in \mathbb{N}, |f_j| \leq C \frac{b^j}{\Gamma(j/p + 1)} \right\}. \quad (2.2)$$

The following lemmas are refinements of results previously stated in [18], except that we need here some extra dependency with respect to auxiliary parameters. They will be of crucial importance in order to prove the main results in the next sections.

Lemma 2.2 *Let \mathcal{T} be a set of parameters and $\tau \in \mathcal{T} \mapsto \mathbb{D}(\tau)$ be a differential operator-valued map*

$$\tau \in \mathcal{T} \mapsto \mathbb{D}(\tau) = \sum_{j=0}^{\infty} b_j(\tau) \left(\frac{d}{dW} \right)^j$$

(with $b_j : \mathcal{T} \rightarrow \mathbb{C}$ for $j \in \mathbb{N}$) whose formal symbol

$$\mathbb{F} : (\tau, W) \in \mathcal{T} \times \mathbb{C} \mapsto \sum_{j=0}^{\infty} b_j(\tau) W^j$$

realizes for each $\tau \in \mathcal{T}$ an entire function of W such that

$$\sup_{\tau \in \mathcal{T}, W \in \mathbb{C}} (|\mathbb{F}(\tau, W)| e^{-B|W|^p}) = A < +\infty \quad (2.3)$$

for some $p \geq 1$ and $B \geq 0$. Then $\mathbb{D}(\tau)$ acts as a continuous operator from $A_1(\mathbb{C})$ into itself uniformly with respect to the parameter $\tau \in \mathcal{T}$.

Proof It follows from Lemma 2.1 that the coefficient functions $\tau \mapsto b_j(\tau)$ satisfy then uniform estimates

$$\forall j \in \mathbb{N}, \forall \tau \in \mathcal{T}, |b_j(\tau)| \leq C \frac{b^j}{\Gamma(j/p + 1)}$$

for some positive constants $C = C(\mathbb{D})$ and $b = b(\mathbb{D})$ depending only on the finite quantity A in (2.3) and B . Let $f : W \mapsto \sum_{\ell=0}^{\infty} a_{\ell} W^{\ell} \in A_1(\mathbb{C})$. There are then (see again Lemma 2.1) positive constants γ and β such that $|a_{\ell}| \leq (\gamma/\ell!) \beta^{\ell}$ for any $\ell \in \mathbb{N}$. Consider the action of \mathbb{D} on such f . One has (for the moment formally)

$$\forall \tau \in \mathcal{T}, \mathbb{D}(\tau)(f) = \sum_{j=0}^{\infty} b_j(\tau) (d/dW)^j (f) = \sum_{j=0}^{\infty} b_j(\tau) \left(\sum_{\ell=0}^{\infty} \frac{(j+\ell)!}{\ell!} a_{\ell+j} W^{\ell} \right)$$

$$= \sum_{\ell=0}^{\infty} \left(\sum_{j=0}^{\infty} \frac{(j+\ell)!}{\ell!} b_j(\tau) a_{\ell+j} \right) W^\ell \tag{2.4}$$

with

$$\sum_{j=0}^{\infty} \frac{(j+\ell)!}{\ell!} |b_j(\tau)| |a_{\ell+j}| \leq \gamma C \frac{\beta^\ell}{\ell!} \sum_{j=0}^{\infty} \frac{(b\beta)^j}{\Gamma(j/p+1)} = K(b, C, \beta, \gamma) \frac{\beta^\ell}{\ell!}. \tag{2.5}$$

Therefore the formal identity (2.4) is in fact a true one for any $W \in \mathbb{C}$, which shows that $\mathbb{D}(\tau)[f] \in A_1(\mathbb{C})$ for any $\tau \in \mathcal{T}$, with

$$\forall \tau \in \mathcal{T}, \forall W \in \mathbb{C}, \quad |\mathbb{D}(\tau)(f)| \leq K(b, C, \beta, \gamma) e^{\beta|W|}.$$

Let $\mathbf{f} = \{f_N\}_{N \geq 1}$ be a sequence converging towards 0 in $A_1(\mathbb{C})$ which is equivalent to say that $\sup(b_{f_N} + C_{f_N}) < +\infty$ and that \mathbf{f} converges towards 0 in $H(\mathbb{C})$, see Proposition 2.1. Then the sequence $\{\mathbb{D}(\tau)(f_N)\}_{N \geq 1} = \mathbb{D}(\tau)(\mathbf{f})$ is such that

$$\forall N \in \mathbb{N}^*, \forall \tau \in \mathcal{T}, \forall W \in \mathbb{C}, \quad |\mathbb{D}(\tau)(f_N)(W)| \leq A_f e^{B_f|W|}$$

for some positive constants A_f and B_f depending only on \mathbb{D} and \mathbf{f} . Let $\mathfrak{B} > B_f$ and $\varepsilon > 0$. Let $R = R_\varepsilon$ large enough such that

$$\forall \tau \in \mathcal{T}, \forall N \in \mathbb{N}^*, \forall W \in \mathbb{C} \text{ with } |W| > R, \quad |\mathbb{D}(\tau)(f_N)(W)| e^{-\mathfrak{B}|W|} \leq \varepsilon.$$

Since $\mathbb{D}(\tau)(f_N)(W) = \sum_{\ell=0}^{\infty} a_{N,\ell}(\tau) W^\ell$ with $|a_{N,\ell}(\tau)| \leq (C_f/\ell!) b_f^\ell$ for some constants C_f and b_f independent on $\tau \in \mathcal{T}$ and on N (see (2.5)) and the sequence \mathbf{f} converges to 0 in $H(\mathbb{C})$, one can find $N = N_\varepsilon$ such that

$$\forall N \geq N_\varepsilon, \forall \tau \in \mathcal{T}, \forall W \in \mathbb{C} \text{ with } |W| \leq R, \quad |\mathbb{D}(\tau)(f_N)(W)| \leq \varepsilon.$$

Hence the sequence $\mathbb{D}(\tau)(\mathbf{f})$ converges towards 0 in $A_1(\mathbb{C})$, uniformly with respect to the parameter τ . □

Since the next lemma involves as set of parameters \mathcal{T} the set which is now given as split in the form $\mathcal{T} = \mathfrak{T} \times \mathbb{C}_Z$, where \mathbb{C}_Z is already a copy of the complex plane, one needs to duplicate \mathbb{C}_Z into an extra copy of \mathbb{C} denoted as \mathbb{C}_W .

Lemma 2.3 *Let \mathfrak{T} be a set of parameters and $\mathfrak{t} \in \mathfrak{T} \mapsto \mathbb{D}(\mathfrak{t}, Z)$ be a differential operator-valued map*

$$\mathfrak{t} \in \mathfrak{T} \mapsto \mathbb{D}(\mathfrak{t}, Z) = \sum_{j=0}^{\infty} b_j(\mathfrak{t}, Z) \left(\frac{d}{dZ} \right)^j$$

(with $b_j : \mathfrak{T} \times \mathbb{C} \rightarrow \mathbb{C}$, holomorphic in Z for $j \in \mathbb{N}$) such that

$$\forall \varepsilon > 0, \quad \sup_{t \in \mathfrak{T}, (Z, W) \in \mathbb{C}^2} \left(\left(\sum_{j=0}^{\infty} |b_j(t, Z)| |W|^j \right) \exp(-\varepsilon |Z|^{\check{p}} - B |W|^p) \right) = A^{(\varepsilon)} < +\infty \quad (2.6)$$

for some $\check{p} > 1$, $p \geq 1$ and $B \geq 0$. Then $\mathbb{D}(t, Z)$ acts as a continuous operator from $A_1(\mathbb{C})$ into $A_{\check{p},0}(\mathbb{C})$ uniformly with respect to the parameter $t \in \mathfrak{T}$.

Proof The function

$$\mathbb{F} : (t, Z, W) \mapsto \sum_{j=0}^{\infty} \left(\sum_{\kappa=0}^{\infty} b_{j,\kappa}(t) Z^\kappa \right) W^j = \sum_{\kappa=0}^{\infty} Z^\kappa \left(\sum_{j=0}^{\infty} b_{j,\kappa}(t) W^j \right) \quad (2.7)$$

is well defined and depends as an entire function of two variables of the variables Z and W (which also justifies in (2.7) the application of Fubini theorem). Cauchy formulae in $\mathbb{C} \times \mathbb{C}$ show that for any $t \in \mathfrak{T}$, for any $j, \kappa \in \mathbb{N}$,

$$\begin{aligned} |b_{j,\kappa}(t)| &= \frac{1}{4\pi^2} \left| \int_{|Z|=\check{r}, |W|=r} F(t, Z, W) \frac{dZ}{Z^{\kappa+1}} \wedge \frac{dW}{W^{j+1}} \right| \leq A^{(\varepsilon)} \inf_{\check{r}>0} \frac{e^{\varepsilon \check{r}^{\check{p}}}}{r^\kappa} \times \inf_{r>0} \frac{e^{Br^p}}{r^j} \\ &= A^{(\varepsilon)} \left(\frac{1}{\kappa} \right)^{\kappa/\check{p}} \times \left(\frac{1}{j} \right)^{j/p} ((\varepsilon \check{p} e)^{1/\check{p}})^\kappa ((Bpe)^{1/p})^j \\ &\leq C_\eta \frac{1}{\Gamma(\kappa/\check{p} + 1)\Gamma(j/p + 1)} (\eta \check{b})^\kappa b^j \end{aligned} \quad (2.8)$$

for each $\eta > 0$, with constants C_η, \check{b} and b independent on the parameter t . Let now $f = \{f_N\}_{N \geq 1}$ be a sequence of elements in $A_1(\mathbb{C})$ which converges to 0 in $A_1(\mathbb{C})$. All differential operators

$$\mathbb{D}_\kappa(t) := \sum_{j=0}^{\infty} b_{j,\kappa}(t) (d/dW)^j \quad (\kappa \in \mathbb{N})$$

act continuously on $A_1(\mathbb{C})$, as seen in Lemma 2.2. Moreover, one has (plugging in (2.5) the estimates (2.8)) that

$$\begin{aligned} \forall f \in A_1^{\gamma,\beta}(\mathbb{C}), \quad \forall t \in \mathfrak{T}, \quad \forall \kappa \in \mathbb{N}, \quad \forall \ell \in \mathbb{N}, \\ (\mathbb{D}_\kappa(t)(f))_\ell \leq \gamma \tilde{C}_\eta \frac{(\eta \check{b})^\kappa}{\Gamma(\kappa/\check{p} + 1)} E_{1/p,1}(\beta b) \frac{\beta^\ell}{\ell!} \end{aligned}$$

where $E_{1/p,1} : \zeta \in \mathbb{C} \mapsto \sum_{k=0}^{\infty} \zeta^k / \Gamma(k/p + 1)$ is the entire (with order $1/p$ and type 1) Mittag–Leffler function. One has therefore for such $f \in A_1^{\gamma,\beta}(\mathbb{C}_W)$ that

$$\forall t \in \mathfrak{T}, \forall \kappa \in \mathbb{N}, \forall W \in \mathbb{C}, \quad |\mathbb{D}_\kappa(t)(f)(W)| \leq \gamma C_\eta E_{1/p,1}(\beta b) \frac{e^{\beta|W|}}{\Gamma(\kappa/\check{p} + 1)} \tag{2.9}$$

and (taking now $W = Z$)

$$\forall t \in \mathfrak{T}, \forall Z \in \mathbb{C}, \quad \sum_{\kappa=0}^\infty |Z|^\kappa |\mathbb{D}_\kappa(t)(f)(Z)| \leq \gamma C_\eta E_{1/p,1}(\beta b) e^{\beta|Z|} \sum_{\kappa=0}^\infty \frac{(\eta \check{b} |Z|)^\kappa}{\Gamma(\kappa/\check{p} + 1)}. \tag{2.10}$$

Since the Mittag–Leffler function $E_{1/\check{p},1}$ has order $\check{p} > 1$, the estimates (2.10) (uniform in the parameter t as well as on the function $f \in A_1^{\gamma,\beta}(\mathbb{C})$) show that the differential operator acts continuously from $A_1(\mathbb{C})$ into $A_{\check{p},0}(\mathbb{C})$, uniformly with respect to the parameter $t \in \mathfrak{T}$. One just needs to repeat here the end of the proof of Lemma 2.2. \square

We conclude this section by proving a quantitative lemma which reveals to be essential in the sequel. It is a refinement of Lemma 1 in [34].

Lemma 2.4 *Let $a \in \mathbb{C}$ with $\alpha := \max(1, |a|)$ and, for any $z \in \mathbb{C}$,*

$$F_N(z, a) := \left(\cos\left(\frac{z}{N}\right) + i a \sin\left(\frac{z}{N}\right) \right)^N$$

as in (1.1) (with $z, a \in \mathbb{C}$ instead of $x, a \in \mathbb{R}$). For any $N \in \mathbb{N}^$ and any $z \in \mathbb{C}$, one has*

$$\begin{aligned} |F_N(z, a)| &\leq \exp(|a||z| + |\operatorname{Im}(z)|) \leq \exp((|a| + 1)|z|) \\ |F_N(z, a) - e^{iaz}| &\leq \frac{2}{3} \frac{|a^2 - 1|}{N} |z|^2 \exp((\alpha + 1)|z|). \end{aligned} \tag{2.11}$$

Proof Let

$$\operatorname{sinc} : z \in \mathbb{C} \mapsto \frac{\sin z}{z} = \int_0^1 t \cos(tz) dt$$

be the sinus cardinal function; it satisfies $|\operatorname{sinc}(z)| \leq e^{|\operatorname{Im}(z)|}$ for any $z \in \mathbb{C}$. One has then the upper uniform estimates

$$\begin{aligned} \forall N \in \mathbb{N}^*, \forall z \in \mathbb{C}, |F_N(z, a)| &= \left| \cos\left(\frac{z}{N}\right) + ia \sin\left(\frac{z}{N}\right) \right|^N \\ &= \left| \cos\left(\frac{z}{N}\right) + i \frac{az}{N} \operatorname{sinc}\left(\frac{z}{N}\right) \right|^N \\ &\leq e^{|\operatorname{Im}(z)|} \left(1 + \frac{|az|}{N}\right)^N \leq \exp(|a||z| + |\operatorname{Im}(z)|) \leq \exp((|a| + 1)|z|), \end{aligned} \tag{2.12}$$

which is the first chain of inequalities in (2.11). For any $N \in \mathbb{N}^*$, one has also

$$\left| \cos\left(\frac{z}{N}\right) - \cos\left(\frac{az}{N}\right) \right| = 2 \left| \sin\left(\frac{(a-1)z}{2N}\right) \sin\left(\frac{(a+1)z}{2N}\right) \right|$$

$$\leq \frac{|a^2 - 1|}{2N^2} |z|^2 \exp\left(\frac{|a - 1| + |a + 1|}{2N} |z|\right) \leq \frac{|a^2 - 1|}{2N^2} |z|^2 \exp\left(\frac{\alpha + 1}{N} |z|\right) \quad (2.13)$$

and

$$\begin{aligned} \left| a \sin\left(\frac{z}{N}\right) - \sin\left(\frac{az}{N}\right) \right| &= \left| \sum_{k=0}^{\infty} \frac{(-1)^k}{(2k + 1)!} (a - a^{2k+1}) \left(\frac{z}{N}\right)^{2k+1} \right| \\ &= \frac{|a^2 - 1|}{N^2} |z|^2 \left| \sum_{k=1}^{\infty} \frac{(-1)^k}{(2k + 1)!} \left(\sum_{\ell=0}^{k-1} a^{2\ell+1} \right) \left(\frac{z}{N}\right)^{2k-1} \right| \\ &\leq \frac{|a^2 - 1|}{2N^2} |z|^2 \sum_{k=1}^{\infty} \frac{\alpha^{2k-1}}{(2k - 1)!(2k + 1)} \left(\frac{|z|}{N}\right)^{2k-1} \\ &\leq \frac{|a^2 - 1|}{6N^2} |z|^2 \sum_{k=1}^{\infty} \frac{1}{(2k - 1)!} \left(\frac{\alpha|z|}{N}\right)^{2k-1} \leq \frac{|a^2 - 1|}{6N^2} |z|^2 \exp\left(\frac{\alpha}{N} |z|\right). \end{aligned} \quad (2.14)$$

It follows from the identity $A^N - B^N = (A - B) \sum_{k=0}^{N-1} A^k B^{N-1-k}$, together with estimates (2.13), (2.14) and (2.12), that for any $N \in \mathbb{N}^*$ and $z \in \mathbb{C}$,

$$\begin{aligned} |F_N(z, a) - e^{iaz}| &= \left| \cos\left(\frac{z}{N}\right) - \cos\left(\frac{az}{N}\right) + i\left(a \sin\left(\frac{z}{N}\right) - \sin\left(\frac{az}{N}\right)\right) \right| \\ &\quad \times \sum_{k=0}^{N-1} |F_N(z, a)|^{k/N} \left| \exp\left(iaz \frac{N-1-k}{N}\right) \right| \\ &\leq \frac{2}{3} \frac{|a^2 - 1|}{N^2} |z|^2 \exp\left(\frac{\alpha + 1}{N} |z|\right) \sum_{k=0}^{N-1} \exp\left(k \left(\frac{|a| + 1}{N}\right) |z| + \frac{N-1-k}{N} |a| |z|\right) \\ &\leq \frac{2}{3} \frac{|a^2 - 1|}{N} |z|^2 \exp((\alpha + 1)|z|). \end{aligned}$$

The second inequality in (2.11) is thus proved. □

One can now state as a consequence of Proposition 2.1 and Lemma 2.4 the following theorem.

Theorem 2.1 *For any $a \in \mathbb{C}$, the sequence $\{z \mapsto F_N(z, a)\}_{N \geq 1}$ converges to $z \mapsto e^{iaz}$ in $A_1(\mathbb{C})$.*

Proof It follows from estimates (2.12) that the sequence $f = \{z \mapsto F_N(z, a)\}_{N \geq 1}$ satisfies the estimates (2.1) with $p = 1$, $B_f = |a| + 1$ and $C_f = 1$. Lemma 2.4 implies on the other hand that the sequence f converges towards $z \mapsto e^{iaz}$ in $H(\mathbb{C})$. The result is then a consequence of Proposition 2.1. □

3 Uniform Convergence of Superscillating Sequences

Let $m \in \mathbb{N}^*$ and $(\mathcal{F}(\mathbb{R}^m, \mathbb{C}))^{\mathbb{N}^*}$ be the family of all sequences $Y = \{x \in \mathbb{R}^m \mapsto Y_N(x)\}_{N \geq 1}$ of complex valued functions defined on \mathbb{R}^m . We first recall in this section the notions of (complex) *generalized Fourier sequence* (CGFS) and (complex) *superscillating sequence* (CSOscS) in $(\mathcal{F}(\mathbb{R}^m, \mathbb{C}))^{\mathbb{N}^*}$. We start first with the case $m = 1$.

Definition 3.1 A sequence $Y \in (\mathcal{F}(\mathbb{R}, \mathbb{C}))^{\mathbb{N}^*}$ is called a complex generalized Fourier sequence if each entry Y_N is, after re-indexation, of the form

$$Y_N : x \in \mathbb{R} \mapsto \sum_{j=0}^N C_j(N) \exp(ik_j(N)x), \tag{3.1}$$

where $C_j(N) \in \mathbb{C}$ and $k_j(N) \in \mathbb{R}$ for any $N \in \mathbb{N}^*$ and $j = 0, \dots, N$.

Example 3.1

1. If $f \in L^1(\mathbb{T}, \mathbb{C})$, where $\mathbb{T} = \mathbb{R}/(2\pi\mathbb{Z})$, is any subsequence of the Fourier (resp. Fourier-Fejér) sequences $\{x \mapsto S_N(x)\}_{N \geq 1}$ (resp. $\{x \mapsto F_N(x)\}_{N \geq 1}$), where

$$S_N(x) = \sum_{j=0}^{2N} \left(\int_{\mathbb{T}} f(\theta) e^{-i(j-N)\theta} \frac{d\theta}{(2\pi)} \right) e^{i(j-N)x}$$

$$F_N(x) = \sum_{j=0}^{2N} \left(1 - \frac{|j-N|}{N} \right) \left(\int_{\mathbb{T}} f(\theta) e^{-i(j-N)\theta} \frac{d\theta}{(2\pi)} \right) e^{i(j-N)x},$$

then it realizes, after re-indexation, an archetypical example of a complex generalized Fourier sequence in $(\mathcal{F}(\mathbb{R}, \mathbb{C}))^{\mathbb{N}^*}$. This fact justifies the terminology.

2. When $m = 1$ and $a \in \mathbb{R}$, the sequence $\{x \mapsto F_N(x, a)\}_{N \geq 1}$ is also an example of a complex generalized Fourier sequence in $(\mathcal{F}(\mathbb{R}, \mathbb{C}))^{\mathbb{N}^*}$. In this case, note that $C_j(N) = C_j(N, a) \in \mathbb{R}$ for any $j \in \mathbb{N}$.
3. Let $P = \sum_{\kappa \in \mathbb{Z}^*} \gamma_\kappa X^\kappa \in \mathbb{C}[X, X^{-1}]$ be a Laurent polynomial and $L(P)$ the diameter of its support. Any sequence $\{x \mapsto Y_N(x)\}_{N \geq 1}$ such that

$$Y_N(x) = \sum_{j=0}^N C_j(N) P(e^{ik_j(N)x}) = \sum_{j=0}^N \sum_{\kappa \in \mathbb{Z}^*} \lambda_\kappa C_j(N) e^{i\kappa k_j(N)x}$$

$$= \sum_{j=0}^{L(P)N} \tilde{C}_j(N) e^{i\tilde{k}_j(N)x}$$

is after re-indexation a complex generalized Fourier sequence in $(\mathcal{F}(\mathbb{R}, \mathbb{C}))^{\mathbb{N}^*}$.

Definition 3.2 A complex generalized Fourier sequence $\{x \mapsto Y_N(x)\}_{N \geq 1}$ in $(\mathcal{F}(\mathbb{R}, \mathbb{C}))^{\mathbb{N}^*}$ is called a complex superscillating sequence if

- each entry Y_N is of the form (3.1) with $|k_j(N)| \leq 1$ for any $j \in \mathbb{N}$ such that $0 \leq j \leq N$;
- there exists an open subset $U^{\text{soosc}} \subseteq \mathbb{R}$ which is called a superoscillation domain such that $\{x \mapsto Y_N(x)\}_{N \geq 1}$ converges uniformly on any compact subset of U^{soosc} to the restriction to U^{soosc} of a trigonometric polynomial function

$$Y_\infty : x \mapsto P_\infty(e^{ik(\infty)x})$$

where $P_\infty \in \mathbb{C}[X, X^{-1}]$ is a Laurent polynomial with no constant term and $k(\infty) \in \mathbb{R} \setminus [-1, 1]$.

Remark 3.1 If Y is a superoscillating sequence in the sense of Definition 3.2, it is Y_∞ -superoscillating in the sense of Definition 1.1 in [33], with *superoscillation set* any segment $[a, b]$ such that $b - a > 0$ is included in the superoscillation domain U^{soosc} .

Example 3.2 1. Any subsequence of the Fourier (resp. Fourier-Fejér) sequences $\{x \mapsto S_N(x)\}_{N \geq 1}$ (resp. $\{x \mapsto F_N(x)\}_{N \geq 1}$) introduced in Example 3.1 (1) fails to be superoscillating since the condition $|k_j(N)| \leq 1$ is not fulfilled.
 2. If $a \in \mathbb{R} \setminus [-1, 1]$, the sequence $\{x \mapsto F_N(x, a)\}_{N \geq 1}$ is a superoscillating sequence in $(\mathcal{F}(\mathbb{R}, \mathbb{C}))^{\mathbb{N}^*}$ with superoscillation domain equal to \mathbb{R} , with $Y_\infty : x \in \mathbb{R} \mapsto e^{iax}$. This follows from Lemma 2.4 (namely from the inequalities (2.11) for $a \in \mathbb{R}$ and $x \in \mathbb{R}$). This is the model that inspired us originally and that we will generalize in this paper.

Inspired by physical considerations which we will discuss later on, we extend as follows Definitions 3.1 and 3.2 to the higher dimensional setting where $m > 1$. The model we will use in order to extend Definition 3.1 will be the one in Example 3.1 (3).

Definition 3.3 A sequence $Y \in (\mathcal{F}(\mathbb{R}^m, \mathbb{C}))^{\mathbb{N}^*}$ is called a complex generalized Fourier sequence if, after re-indexation, each entry Y_N is of the form

$$Y_N : x = (x_1, \dots, x_m) \in \mathbb{R}^m \mapsto \sum_{j=0}^N C_j(N) P(e^{ix_1 k_{j,1}(N)}, \dots, e^{ix_m k_{j,m}(N)}) \quad (3.2)$$

where $P \in \mathbb{C}[X_1, \dots, X_m] \in \mathbb{C}[X_1^{\pm 1}, \dots, X_m^{\pm 1}]$ is a Laurent polynomial (independent of N), $C_j(N) \in \mathbb{C}$ and $N \mapsto k_j(N)$ is a map from \mathbb{N}^* to \mathbb{R}^m for any $N \in \mathbb{N}^*$ and $j = 0, \dots, N$.

Example 3.3 Let t, x be two real variables, $C_j(N) \in \mathbb{C}$, $\kappa_j(N) \in \mathbb{R}$, $k_j(N) \in \mathbb{R}$ for any $N \in \mathbb{N}^*$ and $0 \leq j \leq N$. Then

$$\left\{ x \mapsto \sum_{j=0}^N C_j(N) e^{i\kappa_j(N)t} e^{ik_j(N)x} \right\}_{N \geq 1}$$

is a complex generalized Fourier sequence in the two real variables t, x , the polynomial $P \in \mathbb{C}[T, X]$ being here $P(T, X) = TX$.

Definition 3.2 extends to the multivariate case as follows.

Definition 3.4 A complex generalized Fourier sequence $\{x \mapsto Y_N(x)\}_{N \geq 1}$ in $(\mathcal{F}(\mathbb{R}^m, \mathbb{C}))^{\mathbb{N}^*}$ is called a complex superscillating sequence if

- each entry Y_N is of the form (3.2) with additionally $|k_{j,\ell}(N)| \leq 1$ for any $j \in \mathbb{N}$ such that $0 \leq j \leq N$ and $\ell = 1, \dots, m$;
- there exists an open subset $U^{\text{osoc}} \subseteq \mathbb{R}^m$ which is called a superscillation domain such that $\{x \mapsto Y_N(x)\}_{N \geq 1}$ converges uniformly on any compact subset of U^{osoc} to the restriction to U^{osoc} of a trigonometric polynomial function

$$Y_\infty : x \mapsto P_\infty(e^{ik_1(\infty)x_1}, \dots, e^{ik_m(\infty)x_m})$$

where $P_\infty \in \mathbb{C}[X_1^{\pm 1}, \dots, X_m^{\pm 1}]$ is a Laurent polynomial with no constant term and $k_j(\infty) \in (\mathbb{R} \setminus [-1, 1])^m$.

In order to illustrate Definition 3.4 with an example which is derived from Example 3.2 (2), consider, for $p \in \mathbb{N}$ and $a \in \mathbb{R} \setminus [-1, 1]$, the complex generalized Fourier sequence in two real variables t, x

$$\left\{ \psi_{p,N}(\cdot, \cdot, a) : (t, x) \in \mathbb{R}^2 \mapsto \sum_{j=0}^N C_j(N, a) e^{i(1-2j/N)^p t} e^{i(1-2j/N)x} \right\}_{N \geq 1} \quad (3.3)$$

(see Example 3.3). An immediate computation shows that for any $(t, x) \in \mathbb{R}^2$,

$$\begin{aligned} \frac{\partial}{\partial t} (\psi_{p,N}(t, x, a)) &= i \sum_{j=0}^N C_j(N, a) (1 - 2j/N)^p e^{i(1-2j/N)^p t} e^{i(1-2j/N)x} \\ \frac{\partial^p}{\partial x^p} (\psi_{p,N}(t, x, a)) &= i^p \sum_{j=0}^N C_j(N, a) (1 - 2j/N)^p e^{i(1-2j/N)^p t} e^{i(1-2j/N)x}, \end{aligned}$$

which shows that $(t, x) \in \mathbb{R}^2 \mapsto \psi_{p,N}(t, x, a)$ is the (unique) global solution of the Cauchy–Kowalevski problem

$$\left(i^{p-1} \frac{\partial}{\partial t} - \frac{\partial^p}{\partial x^p} \right) (\psi) \equiv 0, \quad [\psi(t, x)]|_{t=0} = F_N(x, a). \quad (3.4)$$

One can extend analytically $\psi_{p,N}(\cdot, \cdot, a)$ as a function from $\mathbb{R} \times \mathbb{C}$ to \mathbb{C} , such that one has formally

$$\begin{aligned} \psi_{p,N}(t, z, a) &= \sum_{j=0}^N C_j(N, a) \left(\sum_{\ell=0}^{\infty} \frac{i^{\ell(1-p)} t^\ell}{\ell!} (i(1 - 2j/N))^{p\ell} \right) e^{i(1-2j/N)z} \\ &= \left(\sum_{\ell=0}^{\infty} \frac{i^{\ell(1-p)} t^\ell}{\ell!} D^{p\ell} \right) (F_N(\cdot, a))(z) = \mathbb{D}_p(t)(F_N(\cdot, a))(z). \end{aligned} \quad (3.5)$$

One can prove here the following result.

Theorem 3.1 *The operator $\mathbb{D}_p(t)$ acts continuously from $A_1(\mathbb{C})$ into itself. The generalized Fourier sequence (3.3) is superoscillating with \mathbb{R}^2 as superoscillation domain and limit function*

$$Y_\infty : (t, x) \mapsto e^{ita^p} e^{iax},$$

($P_\infty(T, X) = TX, k_1(\infty) = a^p, k_2(\infty) = a$) uniformly on any compact in \mathbb{R}^2 . For any $(\mu, \nu) \in \mathbb{N}^2$, the sequence of functions

$$\begin{aligned} \frac{\partial^{\mu+\nu}}{\partial t^\mu \partial x^\nu} (\psi_{p,N}(t, x, a)) &= i^{-\mu(1-p)} \frac{\partial^{p\mu+\nu}}{\partial x^{p\mu+\nu}} (\psi_{p,N}(t, x, a)) \\ &= i^{-\mu(1-p)} ((d/dW)^{p\mu+\nu} \odot \mathbb{D}_p(t))(F_N(\cdot, a)(x)) \quad (N \in \mathbb{N}^*) \end{aligned} \tag{3.6}$$

converges uniformly on any compact in \mathbb{R}^2 to the function

$$(t, x) \in \mathbb{R}^2 \mapsto ((d/dW)^{p\mu+\nu} \odot \mathbb{D}_p(t))(e^{ia^p t} e^{ia(\cdot)})(x).$$

Proof The first assertion follows from Lemma 2.2 with $\mathbb{R}_t = \mathcal{I}$ as set of parameters and $p \geq 1$ as order of the symbol of the differential operator $\mathbb{D}_p(t)$ as a differential operator in W . Since $\{z \mapsto F_N(z, a)\}_{N \geq 1}$ converges to $z \mapsto e^{iaz}$ in $A_1(\mathbb{C})$ (see theorem 2.1), the sequence $\{z \mapsto \mathbb{D}_p(t)(F_N(\cdot, a))(z)\}_{N \geq 1}$ converges towards $z \mapsto \mathbb{D}_p(t)(e^{ia(\cdot)})(z)$ locally uniformly with respect to $t \in \mathbb{R}$. One can check that $\mathbb{D}_p(t)(e^{ia(\cdot)})(z) = e^{ia^p t} e^{iaz}$ thanks to an immediate computation. Since $(1 - 2j/N)^p$ and $(1 - 2j/N)$ lie in $[-1, 1]$ for any $j \in \{0, \dots, N\}$ and $a \in \mathbb{R} \setminus [-1, 1]$, the generalized Fourier sequence (3.3) is superoscillating with $P_\infty(T, X) = TX, k_1(\infty) = a^p$ and $k_2(\infty) = a$, the superoscillation domain being here \mathbb{R}^2 . The expressions of the partial derivatives in t in terms of the partial derivatives in x in (3.6) follow from the fact that $\psi_{p,N}(\cdot, \cdot, a)$ satisfies the partial differential equation in the Cauchy–Kowalevski problem (3.4). The last assertion in the theorem results from the continuity of the differentiation d/dz as an operator from $A_1(\mathbb{C})$ into itself. \square

Let now $P \in \mathbb{R}[X]$ be an even polynomial $P(X) = \gamma_0 + \gamma_1 X^2 + \dots + \gamma_{2d'} X^{2d'}$ and $a \in \mathbb{R} \setminus [-1, 1]$. Consider in this case the generalized Fourier sequence

$$\left\{ \psi_{P,N}(\cdot, \cdot, a) : (t, x) \in \mathbb{R}^2 \mapsto \sum_{j=0}^N C_j(N, a) e^{iP(1-2j/N)t} e^{i(1-2j/N)x} \right\}_{N \geq 1}. \tag{3.7}$$

As in the previous case, an easy computation shows that the function $\psi_{P,N}(\cdot, \cdot, a)$ is the unique global solution (on the whole space \mathbb{R}^2) of the Cauchy–Kowalevski problem

$$\left(i \frac{\partial}{\partial t} - \check{P} \left(\frac{\partial}{\partial x} \right) \right) (\psi) \equiv 0, \quad [\psi(t, x)]|_{t=0} = F_N(x, a) \tag{3.8}$$

where $\check{P} = \sum_{\kappa'=0}^{d'} (-1)^{\kappa'+1} \gamma_{2\kappa'} X^{2\kappa'}$, and the partial differential operator is here of Schrödinger type. Let us introduce the differential operator $\mathbb{D}_P(t)$ defined as

$$\mathbb{D}_P(t) = \bigodot_{\kappa'=0}^{d'} \left(\sum_{\ell=0}^{\infty} \frac{(i^{1-2\kappa'} t \gamma_{2\kappa'})^\ell}{\ell!} (d/dW)^{2\kappa'\ell} \right)$$

with symbol in $A_{2d'}(\mathbb{C}_W)$ (the set of parameters \mathcal{T} being again $\mathcal{T} = \mathbb{R}_t$).

Theorem 3.2 *Let $P \in \mathbb{R}[X]$ be an even polynomial with degree $2d'$. For any $\lambda \in \mathbb{R}$, the Cauchy–Kowalevski problem (of Schrödinger type)*

$$\left(i \frac{\partial}{\partial t} - \check{P} \left(\frac{\partial}{\partial x} \right) \right) (\psi) \equiv 0, \quad [(t, x) \mapsto \psi(t, x)]|_{t=0} = [x \mapsto e^{i\lambda x}] \quad (3.9)$$

admits as unique global solution in \mathbb{R}^2 the function $(t, x) \mapsto \varphi_\lambda(t, x) = e^{itP(\lambda)} e^{i\lambda x}$. One has $\psi_{P,N}(\cdot, \cdot, \lambda) = \sum_{j=0}^N C_j(N, \lambda) \varphi_{1-2j/N}$ and the sequence $\{(t, x) \mapsto \psi_{P,N}(t, x, \lambda)\}_{N \geq 1}$ converges uniformly on any compact set in \mathbb{R}^2 to $(t, x) \mapsto e^{itP(\lambda)} e^{i\lambda x}$. For any $(\mu, \nu) \in \mathbb{N}^2$, the sequence of functions

$$\begin{aligned} \frac{\partial^{\mu+\nu}}{\partial t^\mu \partial x^\nu} (\psi_{P,N}(t, x, \lambda)) &= (-i)^\mu \left(\check{P} \left(\frac{\partial}{\partial x} \right) \right)^{\odot \mu} \odot \left(\frac{\partial}{\partial x} \right)^{\odot \nu} (\psi_{P,N}(t, x, \lambda)) \\ &= (-i)^\mu \left((\check{P}(d/dW))^{\odot \mu} \odot (d/dW)^\nu \odot \mathbb{D}_P(t) \right) (F_N(\cdot, \lambda))(x) \quad (N \in \mathbb{N}^*) \end{aligned} \quad (3.10)$$

converges uniformly on any compact in \mathbb{R}^2 to the function

$$(t, x) \in \mathbb{R}^2 \mapsto (-i)^\mu \left((\check{P}(d/dW))^{\odot \mu} \odot (d/dW)^\nu \odot \mathbb{D}_P(t) \right) (e^{iP(\lambda)t} e^{i\lambda(\cdot)})(x).$$

Proof One has

$$\left(i \frac{\partial}{\partial t} - \check{P} \left(\frac{\partial}{\partial x} \right) \right) (\varphi_\lambda) = (-P(\lambda) + P(\lambda)) e^{itP(\lambda)} e^{i\lambda x} \equiv 0$$

and $\varphi_\lambda(0, x) = e^{i\lambda x}$ for all $x \in \mathbb{R}$. It follows from Lemma 2.2 that the operator $\mathbb{D}_P(t)$ acts continuously on $A_1(\mathbb{C})$, locally uniformly with respect to the parameter $t \in \mathbb{R}$. Since the sequence $\{z \in \mathbb{C} \mapsto F_N(\cdot, \lambda)\}_{N \geq 1}$ converges to $z \mapsto e^{i\lambda z}$ in $A_1(\mathbb{C})$, the sequence $\{z \in \mathbb{C} \mapsto \mathbb{D}_P(t)(F_N(\cdot, \lambda))(z)\}_{N \geq 1}$ converges to $z \mapsto \mathbb{D}_P(t)(e^{i\lambda(\cdot)})(z) = e^{itP(\lambda)} e^{i\lambda z}$ in $A_1(\mathbb{C})$ locally uniformly with respect to the parameter $t \in \mathbb{R}$. The first equality in (3.10) follows from the fact that $\psi_{P,N}(\cdot, \cdot, \lambda)$ is solution of the Cauchy–Kowalevski problem (3.8). The final assertion follows from the continuity of $d/dz : A_1(\mathbb{C}) \rightarrow A_1(\mathbb{C})$. \square

One can even drop the hypothesis about P and take $P = \sum_{\kappa=0}^d \gamma_\kappa X^\kappa$ as polynomial of degree d in $\mathbb{C}[X]$ with associate polynomial $\check{P} = \sum_{\kappa=0}^d (-i)^{\kappa+1} \gamma_\kappa X^\kappa$. The Cauchy–Kowalevski problem (3.8) is not anymore of the Schrödinger type (since

$\check{P} \notin \mathbb{R}[X]$ in general), which makes the only difference with the case previously studied. Nevertheless, one can state exactly the same result, with this time

$$\mathbb{D}_P(t) = \bigodot_{\kappa=0}^d \left(\sum_{\ell=0}^{\infty} \frac{(i^{1-\kappa} t \gamma_{\kappa})^{\ell}}{\ell!} (d/dW)^{\kappa \ell} \right).$$

Theorem 3.3 *Let $P \in \mathbb{C}[X]$ be a polynomial of degree d . All the assertions in Theorem 3.2 are valid, except that (3.8) is not anymore a Cauchy–Kowalevski problem of the Schrödinger type. When $a \in \mathbb{R} \setminus [-1, 1]$, the generalized Fourier sequence $\{x \mapsto \psi_{B,N}(t, x, a)\}_{N \geq 1}$ is superoscillating for any $t \in \mathbb{R}$. Moreover, given such a and $P \in \mathbb{R}[X]$ such that $\sup_{[-1,1]} |P| \leq 1 < |P(a)|$, the generalized Fourier sequence*

$$\left\{ (t, x) \mapsto \psi_{P,N}(t, x, a) = \sum_{j=0}^N C_j(N, a) \varphi_{1-2j/N}(t, x) \right\}_{N \geq 1}$$

is superoscillating as a generalized Fourier sequence in two variables (t, x) , with \mathbb{R}^2 as domain of superoscillation.

Proof The proof follows that of Theorem 3.2. The sequence $\{x \mapsto \psi_{B,N}(t, x, a)\}_{N \geq 1}$ is superoscillating for any $t \in \mathbb{R}$ since it converges on any compact of \mathbb{R}_x (locally uniformly in t) to $x \mapsto e^{itP(a)} e^{iax}$. As for the last assertion, to define Y_{∞} one takes $P_{\infty}(T, X) = TX, \kappa(\infty) = P(a)$ and $k(\infty) = a$ in Definition 3.4. \square

Let now $E(X) = \sum_{\kappa=0}^{\infty} \gamma_{\kappa} X^{\kappa} \in \mathbb{C}[[X]]$ be a power series with radius of convergence $\rho \in]0, +\infty]$, together with the convolution operator

$$\mathbb{D}_E(t) := \lim_{d \rightarrow +\infty} \bigodot_{\kappa=0}^d \left(\sum_{\ell=0}^{\infty} \frac{(i^{1-\kappa} t \gamma_{\kappa})^{\ell}}{\ell!} (d/dW)^{\kappa \ell} \right)$$

with formal symbol

$$\mathbb{F}_E(t) : W \mapsto \exp \left(it \sum_{\kappa=0}^{\infty} i^{1-\kappa} \gamma_{\kappa} W^{\kappa} \right).$$

Since F and $\sum_{\kappa=0}^{\infty} i^{1-\kappa} \gamma_{\kappa} X^{\kappa}$ share the same radius of convergence $\rho > 0$, $\mathbb{F}_E(t)$ realizes, for each $t \in \mathbb{R}$ an holomorphic function in $D(0, \rho) \subset \mathbb{C}_W$ (with Taylor series about 0 depending on $t \in \mathbb{R}$). More precisely, one has

$$\forall t, W \in \mathbb{R} \times D(0, \rho), \quad \mathbb{F}_E(t)(W) = \sum_{j=0}^{\infty} \left(\sum_{\kappa=0}^{\infty} b_{j,\kappa} t^{\kappa} \right) W^j = \sum_{j=0}^{\infty} b_j(t) W^j,$$

where, for $R > 0$, the radius of convergence of the power series $\sum_{j \geq 0} \left(\sum_{\kappa \geq 0} |\beta_{j,\kappa}| R^\kappa \right) X^j$ is at least equal to ρ .

For any $\lambda \in]-\rho, \rho[$ and $z \in \mathbb{C}$, one has formally

$$\begin{aligned} e^{itE(\lambda)} e^{iz\lambda} &= \lim_{d \rightarrow +\infty} \prod_{\kappa=0}^d \left(\sum_{\ell=0}^{\infty} \frac{(i^{1-\kappa} t \gamma_\kappa)^\ell}{\ell!} (i\lambda)^{\kappa\ell} \right) e^{i\lambda z} \\ &= \lim_{d \rightarrow +\infty} \prod_{\kappa=0}^d \left(\sum_{\ell=0}^{\infty} \frac{(i^{1-\kappa} t \gamma_\kappa)^\ell}{\ell!} (d/dW)^{\kappa\ell} \right) (e^{i\lambda(\cdot)})(z) = \mathbb{D}_E(t)(e^{i\lambda(\cdot)})(z). \end{aligned} \tag{3.11}$$

One requires the following lemma in order to justify the formal relations (3.11).

Lemma 3.1 *When $\rho = +\infty$, the convolution operator $\mathbb{D}_E(t)$ acts continuously locally uniformly with respect to $t \in \mathbb{R}$ from $A_1(\mathbb{C})$ into itself. When $\rho \in]0, +\infty[$ it acts continuously locally uniformly with respect to $t \in \mathbb{R}$ from the space*

$$\left\{ f \in A_1(\mathbb{C}) ; \forall \varepsilon > 0, \exists C_\varepsilon > 0 \text{ such that } |f(W)| \leq C_\varepsilon e^{(\rho-\varepsilon)|W|} \right\} = \lim_{\leftarrow} A_1^{B_{\rho,n}}(\mathbb{C})$$

(where $\{B_{\rho,n}\}_{n \geq 1}$ is a strictly increasing sequence converging to ρ) into itself.

Proof Suppose first that $\rho = +\infty$. Let $R > 0$ and $K \subset [-R, R] \subset \mathbb{R}_t$ be a compact set. One recalls here that the radius of convergence of the power series $\sum_{j \geq 0} \left(\sum_{\kappa \geq 0} |\beta_{j,\kappa}| R^\kappa \right) X^j$ equals $+\infty$. Let $\gamma > 0, \beta > 0$ and $f = \sum_{\ell \geq 0} f_\ell W^\ell \in A_1^{\gamma,\beta}(\mathbb{C})$. One can check as in the proof of Lemma 2.2 (compare to (2.5)) that, for any $t \in K$ and $j \in \mathbb{N}$,

$$\sum_{j=0}^{\infty} \frac{(j+\ell)!}{\ell!} |b_j(t)| |f_{\ell+j}| \leq \gamma \frac{\beta^\ell}{\ell!} \sum_{j=0}^{\infty} \left(\sum_{\kappa=0}^{\infty} |b_{j,\kappa}| R^\kappa \right) \beta^j = K_{\mathbb{D}_E}(\beta, \gamma) \frac{\beta^\ell}{\ell!}.$$

This is indeed enough to conclude as in the proof of Lemma 2.2 that $\mathbb{D}_E(t)$ acts continuously locally uniformly in t from $A_1(\mathbb{C})$ into itself.

Consider now the case where $\rho \in]0, +\infty[$. For any $R > 0$, the radius of convergence of the power series $\sum_{j \geq 0} \left(\sum_{\kappa \geq 0} |\beta_{j,\kappa}| R^\kappa \right) X^j$ is now at least equal to ρ . Repeating the preceding argument (but taking now $\beta \leq \rho - \varepsilon$ for some $\varepsilon > 0$ arbitrary small), one concludes that $\mathbb{D}_E(t)$ acts continuously locally uniformly in t from $\lim_{\leftarrow} A_1^{B_{\rho,n}}(\mathbb{C})$ into itself. □

We can now state the last result of this section.

Theorem 3.4 *Let $E = \sum_{\kappa=0}^{\infty} \gamma_\kappa X^\kappa \in \mathbb{C}[[X]]$ be a power series with radius of convergence $\rho \in]2, +\infty[$. Then $\check{E} := \sum_{\kappa=0}^{\infty} (-i)^{\kappa+1} \gamma_\kappa D^\kappa$ acts continuously from*

$\varprojlim A_1^{B_{\rho,n}}(\mathbb{C})$ into itself. For any $t \in \mathbb{R}$ and $a \in \mathbb{R}$ with $1 < |a| < \rho - 1$, the generalized Fourier sequence

$$\left\{ x \in \mathbb{R} \mapsto \psi_{E,N}(t, x, a) = \sum_{j=0}^N C_j(N, a) e^{iE(1-2j/N)t} e^{ix(1-2j/N)} \right\}_{N \geq 1}$$

is superoscillating. Moreover, for any such a and $(\mu, \nu) \in \mathbb{N}^2$, the sequence of functions

$$\begin{aligned} \frac{\partial^{\mu+\nu}}{\partial t^\mu \partial x^\nu} (\psi_{E,N}(t, x, a)) &= (-i)^\mu \left(\check{E}^{\odot \mu} \odot (d/dW)^\nu \right) (\psi_{P,N}(t, \cdot, a))(x) \\ &= (-i)^\mu \left(\check{E}^{\odot \mu} \odot (d/dW)^\nu \odot \mathbb{D}_E(t) \right) (F_N(\cdot, a))(x) \end{aligned} \tag{3.12}$$

converges then uniformly on any compact in $\mathbb{R}_{t,x}^2$ to the fonction

$$(t, x) \mapsto (-i)^\mu \left(\check{E}^{\odot \mu} \odot (d/dW)^\nu \odot \mathbb{D}_E(t) \right) (e^{itE(a)} e^{ia(\cdot)})(x).$$

Proof The fact that \check{E} acts continuously from $\varprojlim A_1^{B_{\rho,n}}(\mathbb{C}_W)$ into itself follows from Lemma 3.1, considering just \check{E} (independent of the parameter t) instead of $\mathbb{D}_E(t)$. For any $\lambda \in \mathbb{R}$ with $|\lambda| < \rho$, the operator \check{E} then acts on $e^{i\lambda(\cdot)}$ and it is immediate to check that for any $t \in \mathbb{R}$

$$\forall x \in \mathbb{R}, \left[\left(i \frac{\partial}{\partial t} - \check{E} \right) (e^{itE(\lambda)} e^{i\lambda W}) \right]_{W=x} = 0; \tag{3.13}$$

moreover $[(t, x) \mapsto e^{itE(\lambda)} e^{i\lambda x}]_{t=0}$ is $x \mapsto e^{i\lambda x}$. Therefore, for any $a \in \mathbb{R}$ and $N \in \mathbb{N}^*$, one has by linearity (since $\rho > 1$)

$$\forall (t, x) \in \mathbb{R}^2, \left[\left(i \frac{\partial}{\partial t} - \check{E} \right) \left(\sum_{j=0}^N C_j(N, a) e^{iE(1-2j/N)t} e^{i(1-2j/N)W} \right) \right]_{W=x} = 0. \tag{3.14}$$

Lemma 3.1 (applied this time with $\mathbb{D}_E(t)$), combined with Theorem 2.1 and the estimates in the first line of (2.11) in Lemma 2.4, imply that as soon as one has $|a| < \rho - 1$ the sequence $\{z \in \mathbb{C} \mapsto \mathbb{D}_E(t)(F(\cdot, a))(z)\}_{N \geq 1}$ converges (locally uniformly with respect to the parameter t) to $z \mapsto e^{iaz}$ in $A_1(\mathbb{C})$. The last assertion in the particular case $\mu = \nu = 0$ follows. The first equality in (3.12) comes from the identity (3.14), while the second one comes from (3.11) (as justified by Lemma 3.1). The last assertion of the theorem when μ, ν are arbitrary is then a consequence of the continuity of d/dz from $\varprojlim A_1^{B_{\rho,n}}(\mathbb{C})$ into itself. The superoscillating character of the sequence $\{\psi_{P,N}(t, \cdot, a)\}_{N \geq 1}$ follows from Definition 3.2. \square

Remark 3.2 When $E \in \mathbb{R}[[X]]$, $1 < |a| < \rho - 1$ and $\sup_{[-1,1]}(E) \leq 1 < |E(a)|$, the generalized Fourier sequence

$$\left\{ (t, x) \in \mathbb{R}^2 \mapsto \psi_{E,N}(t, x, a) = \sum_{j=0}^N C_j(N, a) e^{iE(1-2j/N)t} e^{ix(1-2j/N)} \right\}_{N \geq 1}$$

is also superscillating, this time according to Definition 3.4 (with $P_\infty(T, X) = TX$, $\kappa(\infty) = E(a)$ and $k(\infty) = a$).

4 Regularization of Formal Fresnel-Type Integrals

In order to settle from the mathematical point of view the approach to non-absolutely convergent integrals on the half-line \mathbb{R}^{+*} or the whole real line \mathbb{R} through the so-called principle of *regularization* that we will invoke in the remaining Sects. 5 and 6 (with respect to supershift considerations related to Schrödinger equations with specific potentials), we need to explain what regularization of formal Fresnel-type integrals on \mathbb{R}^{+*} or \mathbb{R} means.

Suppose that \mathfrak{T} is a set of parameters. Let $G : (t, Z) \in \mathfrak{T} \times \mathbb{C} \mapsto G(t, Z)$ be a function which is entire as a function of Z for each $t \in \mathfrak{T}$ fixed. Let also ϕ be a non-vanishing real function on \mathfrak{T} that will play the role of a *phase function*. Let finally χ be a real number such that $\chi > -1$. In order to give a meaning to the formal integral

$$\int_0^\infty (x')^\chi e^{-i\phi(t)(x')^2} G(t, x') dx' \quad (\chi > -1) \tag{4.1}$$

we distinguish the cases where $\phi(t) > 0$ and $\phi(t) < 0$. In the first case ($\phi(t) > 0$), we rewrite this (for the moment formal) expression (4.1) as

$$\begin{aligned} & \int_0^\infty (x')^\chi e^{-i\phi(t)(x')^2} G(t, x') dx' \\ &= e^{-i(\chi+1)\pi/4} \int_{\mathbb{R}^{+*} e^{i\pi/4}} Z^\chi e^{-\phi(t)Z^2} G(t, e^{-i\pi/4}Z) dZ \\ &= \int_{\mathbb{R}^{+*} e^{i\pi/4}} Z^\chi e^{-\phi(t)Z^2} F_+(t, Z) dZ \end{aligned} \tag{4.2}$$

with $F_+(t, Z) := e^{-i(\chi+1)\pi/4} G(t, e^{-i\pi/4}Z)$ for any $t \in \mathfrak{T}$ and $Z \in \mathbb{C}$. In the second case ($\phi(t) < 0$), we rewrite it as

$$\begin{aligned} & \int_0^\infty (x')^\chi e^{-i\phi(t)(x')^2} G(t, x') dx' \\ &= e^{i(\chi+1)\pi/4} \int_{\mathbb{R}^{+*} e^{-i\pi/4}} Z^\chi e^{\phi(t)Z^2} G(t, e^{i\pi/4}Z) dZ \\ &= \int_{\mathbb{R}^{+*} e^{-i\pi/4}} Z^\chi e^{\phi(t)Z^2} F_-(t, Z) dZ \end{aligned} \tag{4.3}$$

with $F_-(t, Z) := e^{i(\chi+1)\pi/4}G(t, e^{i\pi/4}Z)$ for any $t \in \mathfrak{T}$ and $Z \in \mathbb{C}$. The following elementary lemma will reveal to be essential.

Lemma 4.1 *Let \mathfrak{T}, ϕ, χ as above and $F : \mathfrak{T} \times \mathbb{C} \rightarrow \mathbb{C}$ be a function which is entire in the complex variable Z and satisfies the growth estimates*

$$\forall \varepsilon > 0, \quad \sup_{t \in \mathfrak{T}, Z \in \mathbb{C}} (|F(t, Z)| \exp(-\varepsilon|Z|^{\check{p}})) < +\infty \tag{4.4}$$

for some $\check{p} \in]1, 2]$, that is $F(t, \cdot) \in A_{\check{p},0}(\mathbb{C})$ uniformly in t . Then, for any $u = e^{i\theta}$ with $\theta \in]-\pi/4, \pi/4[$, the integral

$$\int_{\mathbb{R}^{+*u}} Z^\chi e^{-|\phi(t)|Z^2} F(t, Z) dZ \tag{4.5}$$

is absolutely convergent and remains independent of u ; it equals in particular its value for $u = 1$.

Proof The absolute convergence follows from the estimates (4.4), together with the fact that if $u = e^{i\theta}$, $\text{Re}(tu)^2 = t^2 \cos(2\theta) > 0$ for $t > 0$. The fact that the integrals do not depend of u follows from residue theorem (applied on the oriented boundary of conic sectors with apex at the origin). \square

In view of this lemma, the regularization of an integral of the Fresnel-type such as (4.1) consists in the successive two operations:

1. first transform the formal expression (4.1) into one of the representations (4.2) or (4.3) according to $\text{sign}(\phi(t))$;
2. then invoke Lemma 4.1 (provided the required hypothesis are satisfied) and regularize (4.1) as $\int_0^\infty Z^\chi e^{-\phi(t)Z^2} F_+(t, Z) dZ$ when $\phi(t) > 0$ or $\int_0^\infty Z^\chi e^{\phi(t)Z^2} F_-(t, Z) dZ$ when $\phi(t) < 0$.

Remark 4.1 In order to give a meaning (if possible of course) to the formal integral expression

$$\int_{\mathbb{R}} |x'|^\chi e^{-i\phi(t)(x')^2} G(t, x') dx', \tag{4.6}$$

one splits it as

$$\int_0^\infty (x')^\chi e^{-i\phi(t)(x')^2} G(t, x') dx' + \int_0^\infty (x')^\chi e^{-i\phi(t)(x')^2} G(t, -x') dx'$$

and proceed as above for the two formal expressions involved into this formal decomposition.

It is immediate to compare this approach to regularization to the alternative following one.

Proposition 4.1 *Let $G \in A_{2,0}(\mathbb{C})$ and $\chi > -1$. Then, for all $\varpi \in \mathbb{R}^*$*

$$\lim_{\varepsilon \rightarrow 0} \int_0^\infty (x')^\chi e^{i\varpi(x')^2} e^{-\varepsilon(x')^2} G(x') dx'$$

exists and coincides with the integral regularized under the approach described above.

Proof It is enough to prove the result when $\varpi = \pm 1$ since one reduces to one of these two cases up to a homothety on the real half line. One has

$$\begin{aligned} \int_0^\infty (x')^\chi e^{-\varepsilon(x')^2} e^{-i(x')^2} G(x') dx' &= \int_{e^{i\pi/4}\mathbb{R}^{+*}} Z^\chi e^{-(1-i\varepsilon)Z^2} F_+(Z) dZ \\ \int_0^\infty (x')^\chi e^{-\varepsilon(x')^2} e^{i(x')^2} G(x') dx' &= \int_{e^{-i\pi/4}\mathbb{R}^{+*}} Z^\chi e^{-(1+i\varepsilon)Z^2} F_-(Z) dZ, \end{aligned}$$

where $F_+(Z) = e^{-i(1+\chi)\pi/4} F(e^{-i\pi/4}Z)$ and $F_-(Z) = e^{i(1+\chi)\pi/4} F(e^{i\pi/4}Z)$. Let $\rho_\varepsilon = \sqrt{1 + \varepsilon^2}$, and $\xi_\varepsilon = \arg_{[0,\pi/2[} \sqrt{1 + i\varepsilon}$. One has then

$$\begin{aligned} \int_0^\infty (x')^\chi e^{-\varepsilon(x')^2} e^{-i(x')^2} G(x') dx' &= \left(\frac{e^{i\xi_\varepsilon}}{\sqrt{\rho_\varepsilon}}\right)^{1+\chi} \\ &\int_{e^{i(\pi/4-\xi_\varepsilon)}\mathbb{R}^{+*}} Z^\chi e^{-Z^2} F^+(e^{i\xi_\varepsilon} Z/\sqrt{\rho_\varepsilon}) dZ \\ \int_0^\infty (x')^\chi e^{-\varepsilon(x')^2} e^{i(x')^2} G(x') dx' &= \left(\frac{e^{-i\xi_\varepsilon}}{\sqrt{\rho_\varepsilon}}\right)^{1+\chi} \\ &\int_{e^{-i(\pi/4-\xi_\varepsilon)}\mathbb{R}^{+*}} Z^\chi e^{-Z^2} F_-(e^{-i\xi_\varepsilon} Z/\sqrt{\rho_\varepsilon}) dZ. \end{aligned} \tag{4.7}$$

In the two integrals on the right-hand side of the equalities (4.7), the integration contour can be replaced by the half-line \mathbb{R}^{+*} as a consequence of Lemma 4.1. It is then possible to take the limit when ε tends to 0. Lebesgue’s domination theorem then applies and since ρ_ε tends to 1 and ξ_ε to 0, one gets

$$\begin{aligned} \lim_{\varepsilon \rightarrow 0} \int_0^\infty (x')^\chi e^{-\varepsilon(x')^2} e^{-i(x')^2} G(y) dy &= \int_0^\infty Z^\chi e^{-Z^2} F_+(Z) dZ \\ \lim_{\varepsilon \rightarrow 0} \int_0^\infty (x')^\chi e^{-\varepsilon(x')^2} e^{i(x')^2} G(x') dx' &= \int_0^\infty Z^\chi e^{-Z^2} F_-(Z) dZ. \end{aligned}$$

This concludes the proof of the Proposition. □

5 Fresnel-Type Integral Operators

5.1 Continuity on $A_1(\mathbb{C})$ of Fresnel-Type Integral Operators

Let \mathfrak{T} be a set of parameters and $t \in \mathfrak{T} \mapsto \mathbb{D}(t, Z)$ (as in the statement of Lemma 2.3) be a differential operator-valued map

$$t \in \mathfrak{T} \mapsto \mathbb{D}(t, Z) = \sum_{j=0}^{\infty} b_j(t, Z) \left(\frac{d}{dZ}\right)^j$$

(with $b_j : T \times \mathbb{C} \rightarrow \mathbb{C}$, holomorphic in Z for $j \in \mathbb{N}$) such that

$$\forall \varepsilon > 0, \sup_{t \in \mathfrak{T}, (Z, W) \in \mathbb{C}^2} \left(\left(\sum_{j=0}^{\infty} |b_j(t, Z)| |W|^j \right) \exp(-\varepsilon |Z|^{\check{p}} - B |W|^p) \right) = A^{(\varepsilon)} < +\infty \quad (5.1)$$

for some $\check{p} \in]1, 2]$, $p \geq 1$ and $B \geq 0$. Let also ϕ be a non-vanishing real function on \mathfrak{T} and $\chi > -1$. It follows from the estimates (5.1), together with Lemma 4.1, that the regularization approach described in Sect. 4 allows to define the operator

$$t \mapsto \int_0^{\infty} Z^\chi e^{-i\phi(t) Z^2} \sum_{j=0}^{\infty} b_j(t, Z) \left(\frac{d}{dZ}\right)^j (\cdot) dZ. \quad (5.2)$$

One needs to consider for the moment these operators as acting on entire functions of the complex variable Z . For $\alpha \in \mathbb{C}$, let also H_α be the dilation operator $H_\alpha : f \mapsto f(\alpha(\cdot))$ acting on such functions. The symbol \odot still stands for the composition of operators. The discussion is with respect to the sign of $\phi(t)$.

- When $\phi(t) > 0$,

$$\begin{aligned} & \int_0^{\infty} Z^\chi e^{-i\phi(t) Z^2} \left(\sum_{j=0}^{\infty} b_j(t, Z) \left(\frac{d}{dZ}\right)^j (\cdot) \right) dZ \\ &= e^{-i(1+\chi)\pi/4} \int_0^{\infty} y^\chi e^{-\phi(t) y^2} \\ & \left(\sum_{j=0}^{\infty} b_j(t, e^{-i\pi/4} Z) \left(e^{ij\pi/4} \left(\frac{d}{dZ}\right)^j \odot H_{e^{-i\pi/4}} \right) (\cdot) \right) (y) dy. \end{aligned} \quad (5.3)$$

- When $\phi(t) < 0$,

$$\begin{aligned} & \int_0^{\infty} Z^\chi e^{-i\phi(t) Z^2} \left(\sum_{j=0}^{\infty} b_j(t, Z) \left(\frac{d}{dZ}\right)^j (\cdot) \right) dZ \\ &= e^{i(1+\chi)\pi/4} \int_0^{\infty} y^\chi e^{\phi(t) y^2} \\ & \left(\sum_{j=0}^{\infty} b_j(t, e^{i\pi/4} Z) \left(e^{-ij\pi/4} \left(\frac{d}{dZ}\right)^j \odot H_{e^{i\pi/4}} \right) (\cdot) \right) (y) dy. \end{aligned} \quad (5.4)$$

Theorem 5.1 *Suppose that the parameter space \mathfrak{T} is a topological space and that ϕ is continuous. Consider functions $B_j : \mathfrak{T} \times \mathbb{C} \times \mathbb{C} \rightarrow \mathbb{C}$ ($j \in \mathbb{N}$) which are entire in the two complex entries and such that*

$$\forall \varepsilon > 0, \exists A^{(\varepsilon)}, B^{(\varepsilon)} \geq 0 \text{ such that } \forall t \in \mathfrak{T}, \forall Z \in \mathbb{C}, \forall \check{Z} \in \mathbb{C}, \forall W \in \mathbb{C},$$

$$\sum_{j=0}^{\infty} |B_j(t, Z, \check{Z})| |W|^j \leq A^{(\varepsilon)} e^{\varepsilon |Z|^{\check{p}} + B^{(\varepsilon)} |\check{Z}|^{\check{p}} + B |W|^p} \tag{5.5}$$

for some $p \geq 1, \check{p} \in]1, 2]$, and $B \geq 0$. Then the operator

$$\int_0^{\infty} Z^\chi e^{-i\phi(t)Z^2} \left(\sum_{j=0}^{\infty} B_j(t, Z, \check{Z}) \left(\frac{d}{dZ} \right)^j (\cdot) \right) dZ$$

(understood through the process of regularization as described above) acts continuously locally uniformly in t from $A_1(\mathbb{C})$ into $A_{\check{p}}(\mathbb{C})$.

Proof It is enough to consider \mathfrak{T} as a neighborhood of a point t_0 in which $\phi(t) \geq \varepsilon_0 > 0$ (since ϕ is continuous). Let $f = \{Z \mapsto f_N(Z)\}_{N \geq 1}$ be a sequence of elements in $A_1(\mathbb{C})$ that converges towards 0 in $A_1(\mathbb{C})$, which means (see Proposition 2.1) that all f_N belong to some $A_1^{C,b}(\mathbb{C})$ for some constants $C, b > 0$ independent on N (namely $f_N = \sum_{\ell} a_{N,\ell} Z^\ell$ with $|a_{N,\ell}| \leq C b^\ell / \ell!$ for any $\ell \in \mathbb{N}$). It is clear that the operator

$$\sum_{j=0}^{\infty} B_j(t, e^{-i\pi/4} Z, \check{Z}) \left(e^{ij\pi/4} \left(\frac{d}{dZ} \right)^j \odot H_{e^{-i\pi/4}} \right)$$

involved in the integrand of (5.3) is governed by estimates of the form (5.5). It follows then from Lemma 2.3, taking into account estimates (5.5), that for each $N \in \mathbb{N}^*$ the function

$$H(f_N) : (t, Z, \check{Z}) \in \mathfrak{T} \times \mathbb{C} \times \mathbb{C}$$

$$\mapsto \sum_{j=0}^{\infty} B_j(t, e^{-i\pi/4} Z, \check{Z}) \left(e^{ij\pi/4} \left(\frac{d}{dZ} \right)^j \odot H_{e^{-i\pi/4}} \right) (f_N)(Z)$$

is such that for each $\varepsilon > 0$, there exists $\tilde{A}^{(\varepsilon)} \geq 0$ (depending on $\mathfrak{T}, A^{(\varepsilon)}$, the B_j, b and C , but not on the N) such that

$$\forall (t, Z, \check{Z}) \in \mathfrak{T} \times \mathbb{C} \times \mathbb{C}, |H(f_N)(t, Z, \check{Z})| \leq \tilde{A}^{(\varepsilon)} e^{\varepsilon |Z|^{\check{p}} + B^{(\varepsilon)} |\check{Z}|^{\check{p}}}.$$

Take in particular $\varepsilon < \varepsilon_0$. Then the function

$$\check{Z} \in \mathbb{C} \mapsto \int_0^{\infty} y^\chi e^{-\phi(t)y^2} H(f_N)(t, y, \check{Z}) dy$$

is in $A_{\check{p}}(\mathbb{C})$ since it is estimated as

$$\left| \int_0^\infty y^\chi e^{-\phi(t)y^2} H(f_N)(t, y, \check{Z}) dy \right| \leq \tilde{A}^{(\varepsilon)} \left(\int_0^\infty y^\chi e^{-\varepsilon_0 y^2} e^{\varepsilon y^{\check{p}}} dy \right) e^{B^{(\varepsilon)} |\check{Z}|^{\check{p}}} \quad \forall \check{Z} \in \mathbb{C}$$

(remember that $\check{p} \in]1, 2[$). It remains to show that the sequence

$$\left\{ \check{Z} \mapsto \int_0^\infty y^\chi e^{-\phi(t)y^2} H(f_N)(t, y, \check{Z}) dy \right\}_{N \geq 1}$$

converges to 0 in $A_{\check{p}}(\mathbb{C})$. It is enough (see Proposition 2.1) to prove that it converges to 0 uniformly on any closed disk $\overline{D(0, r)}$ in \mathbb{C} . Fix $\varepsilon < \varepsilon_0$ and $\eta > 0$. Choose then $R_\eta \gg 1$ such that

$$\begin{aligned} \forall N \in \mathbb{N}, \quad & \left| \int_{R_\eta}^\infty y^\chi e^{-\phi(t)y^2} H(f_N)(t, y, \check{Z}) dy \right| \\ & \leq \tilde{A}^{(\varepsilon)} \left(\int_0^\infty y^\chi e^{-\varepsilon_0 y^2} e^{\varepsilon y^{\check{p}}} dy \right) e^{B^{(\varepsilon)} |\check{Z}|^{\check{p}}} \leq \eta e^{-B^{(\varepsilon)} r^{\check{p}}} e^{B^{(\varepsilon)} |\check{Z}|^{\check{p}}} \leq \eta. \end{aligned}$$

On $[0, R_\eta]$, one uses the uniform convergence of f towards 0 on any compact set, hence of $H[f]$ on any compact set, to conclude that for $N \geq N_\eta \gg 1$, one has

$$\left| \int_0^{R_\eta} y^\chi e^{-\phi(t)y^2} H(f_N)(t, y, \check{Z}) dy \right| \leq \eta \quad \forall \check{Z} \in \overline{D(0, r)}.$$

Note that our estimates show that the convergence towards 0 in $A_{\check{p}}(\mathbb{C})$ thus obtained is uniform in $t \in \mathfrak{T}$. □

5.2 Superoscillations and Supershifts

Consider the Schrödinger equation

$$i \frac{\partial \psi}{\partial t}(t, x) = (\mathcal{H}(x)(\psi))(t, x) \tag{5.6}$$

where \mathcal{H} denotes the Hamiltonian operator attached to the physical system which is under consideration. Suppose that $\mathbf{Y} = \{x \mapsto Y_N(x)\}_{N \geq 1}$ is a superoscillating sequence. Since

$$\left(i \frac{\partial}{\partial t} - \mathcal{H}(x) \right) (\psi)(t, x) = 0, \quad [\psi(t, x)]_{t=0} = Y(x)$$

is a Cauchy–Kowalevski problem (assuming that x lies in some open set $U \subset \mathbb{R}$ where the Hamiltonian operator is regular), any entry $x \in U \mapsto Y_N(x)$ evolves in a unique way from $t = 0$ towards $t > 0$ as $(t, x) \mapsto \psi_N(t, x)$. Assume in addition that x lies

in the maximal superscillation domain U_{\max}^{suposc} ; the limit function $x \in U \cap U_{\max}^{\text{sosc}} \mapsto Y_{\infty}(x)$ then also evolves from $U \cap U_{\max}^{\text{sosc}}$ into some function $(t, x) \mapsto \psi_{\infty}(t, x)$. A natural question then occurs. As long as the evolution persists (let say for $t \in [0, T]$), is it true that the sequence $\{x \in U \mapsto \psi_N(t, x)\}_{N \geq 1}$ is such that its restriction to $U \cap U_{\max}^{\text{sosc}}$ converges (uniformly on any compact subset of $U \cap U_{\max}^{\text{sosc}}$) to $x \mapsto \psi_{\infty}(t, x)$? If this is the case, one will say that the superscillating character of the sequence Y persists in time through the Schrödinger evolution operator $\partial/\partial t - \mathcal{H}$ which is here considered.

In order to formulate such question in a different way, let us now consider the (t, x) domain $[0, T] \times (U \cap U_{\max}^{\text{sosc}}) = \mathfrak{T} \times (U \cap U_{\max}^{\text{sosc}}) = \mathcal{T}$ as a parameter set and focus on the map $\lambda \in \mathbb{R} \mapsto \varphi_{\lambda}$, where $\varphi_{\lambda} : \mathcal{T} \rightarrow \mathbb{R}$ is evolved to $[0, T] \times U$ (through the Schrödinger operator) from the initial datum $x \in U \mapsto e^{i\lambda x}$, then restricted to the parameter set \mathcal{T} . Previous considerations lead to the following definition, which is inspired by Definition 3.2.

Definition 5.1 Let \mathcal{T} be a locally compact topological space and $\mathcal{F} = \{\varphi_{\lambda}; \lambda \in \mathbb{R}\}$ be a family of \mathbb{C} -valued functions on \mathcal{T} indexed by \mathbb{R} . A sequence $\psi = \{\tau \in \mathcal{T} \mapsto \psi_N(\tau)\}_{N \geq 1}$ of \mathbb{C} -valued functions on \mathcal{T} is called a \mathcal{F} -supershift (or \mathcal{F} admits ψ as a supershift) if

- any entry ψ_N is of the form $\psi_N = \sum_{j=0}^N C_j(N) \varphi_{k_j(N)}$ with $|k_j(N)| \leq 1$ for any $N \in \mathbb{N}^*$ and $0 \leq j \leq N$;
- there exists an open subset \mathcal{U}^{ssh} of \mathcal{T} called a \mathcal{F} -supershift domain such that the sequence $\{\tau \in \mathcal{U}^{\text{ssh}} \mapsto \psi_N(\tau)\}$ converges locally uniformly towards the restriction to \mathcal{U}^{ssh} of a function ψ_{∞} which is a \mathbb{C} -finite linear combination of elements in \mathcal{F} of the form $\varphi_{\nu k(\infty)}$ with $\nu \in \mathbb{Z}^*$, where $k(\infty) \in \mathbb{R} \setminus [-1, 1]$.

Example 5.1 1. If $\mathcal{T} = \mathbb{R}$ and \mathcal{F} denotes the family of characters $x \in \mathbb{R} \mapsto e^{i\lambda x}$ indexed by the dual copy \mathbb{R}_{λ}^* of \mathbb{R}_x , \mathcal{F} -supershifts are the complex superscillating sequences (see Definition 3.2).
 2. Let $a \in \mathbb{R} \setminus [-1, 1]$, $\mathcal{T} = \mathbb{R}_{t,x}^2$ and $\mathcal{F} = \{\varphi_{\lambda}; \lambda \in \mathbb{R}^*\}$ as defined in Theorem 3.2 or Theorem 3.3. For any $a \in \mathbb{R} \setminus [-1, 1]$, the sequence $\{(t, x) \mapsto \psi_{P,N}(t, x, a)\}_{N \geq 1}$ is a \mathcal{F} -supershift which admits $\mathbb{R}_{t,x}^2 = \mathcal{T}$ as \mathcal{F} -supershift domain.

When \mathcal{T} is of the form $[0, T[\times U$, where U is an open subset in \mathbb{R}_x^{m-1} ($m \geq 2$) and $T \in]0, +\infty]$, one can consider as well families $\mathcal{F} = \{\varphi_{\lambda}; \lambda \in \mathbb{R}\}$ of \mathbb{C} -valued distributions in $\mathbb{R} \times U$ with support in $[0, T[\times U$. In order to define in this new context the notion of \mathcal{F} -supershift, one needs to keep the first clause in Definition 5.1 as it is and modify the second clause as follows : “there exists an open subset $\mathcal{U}^{\text{ssh}} = \mathcal{V}^{\text{ssh}} \cap \mathcal{T}$ (where \mathcal{V} is an open subset in $\mathbb{R} \times U$), called a \mathcal{F} -supershift domain, such that the sequence $\{(\psi_N(\tau))_{|\tau \in \mathcal{U}^{\text{ssh}}}\}_{N \geq 1}$ converges weakly in the sense of distributions in \mathcal{V} to the restriction to \mathcal{V} of a distribution $\psi_{\infty} \in \mathcal{D}'(\mathbb{R} \times U, \mathbb{C})$ which is a \mathbb{C} -finite linear combination of elements in \mathcal{F} of the form $\varphi_{\nu k(\infty)}$ with $\nu \in \mathbb{Z}^*$, where $k(\infty) \in \mathbb{R} \setminus [-1, 1]$ ”.

One will need in Sect. 6 a further extension of this concept of \mathcal{F} -supershift to the case where $\mathcal{T} = [0, T[\times U$, $U \subset \mathbb{R}^{m-1}$ with $m \geq 2$ as above, but elements $\varphi_{\lambda} \in \mathcal{F}$

are now *hyperfunctions* in $\mathbb{R} \times U$ with support in \mathcal{T} . The sequence $\{(\psi_N(\tau))|_{\mathcal{V}^{\text{ssh}}}\}_{N \geq 1}$ needs in this case to converge still in the weak sense, but this time *in the sense of hyperfunctions* in \mathcal{V} , towards the restriction to \mathcal{V} of the *hyperfunction* ψ_∞ . The notion of \mathcal{F} -supershift can thus be enlarged to families $\mathcal{F} = \{\varphi_\lambda; \lambda \in \mathbb{R}\}$ of *hyperfunctions* in $\mathbb{R} \times U$ with support in \mathcal{T} .

5.3 The Schrödinger Cauchy Problem with Centrifugal Potential

We will consider in this subsection the case where $U = \{x \in \mathbb{R}; x > 0\}$ and the hamiltonian in (5.6) is $x \in U \mapsto \mathcal{H}(x) = -(\partial^2/\partial x^2)/2 + u/(2x^2)$, where u denotes a real strictly positive physical constant. The corresponding Cauchy–Kowalevski problem (with $[0, +\infty[\times U$ as phase space) is the *Schrödinger Cauchy problem with centrifugal potential*, see [32] for more references. For this Cauchy–Kowalevski problem, the analysis of the evolution $t \mapsto \psi(t, \cdot)$ of the solution $(t, x) \in [0, \infty[\times U \mapsto \psi(t, x)$ from an initial datum $x \in U \mapsto \psi(0, x)$ can be carried through thanks to the explicit form of the Green function $(t, x, x') \mapsto G(t, x, t' = 0, x')$.

Let $\nu = \sqrt{1 + 4u}/2$ and the Bessel function J_ν defined in $\Omega := \mathbb{C} \setminus]-\infty, 0]$ as

$$\begin{aligned} J_\nu : z \in \Omega &\mapsto \left(\frac{z}{2}\right)^\nu \sum_{k=0}^\infty \frac{(-1)^k}{\Gamma(k+1)\Gamma(\nu+k+1)} \left(\frac{z}{2}\right)^{2k} \\ &= \left(\frac{|z|}{2}\right)^\nu e^{i\nu \arg_{]-\pi, \pi[}(z)} \sum_{k=0}^\infty \frac{(-1)^k}{\Gamma(k+1)\Gamma(\nu+k+1)} \left(\frac{z}{2}\right)^{2k} \\ &= \left(\frac{|z|}{2}\right)^\nu e^{i\nu \arg_{]-\pi, \pi[}(z)} E_\nu(z). \end{aligned} \tag{5.7}$$

Then the Green function $(t, x, x') \mapsto G(t, x, 0, x')$ can be explicited in this case as

$$G(t, x, 0, x') = (-i)^{\nu+1} \frac{\sqrt{xx'}}{t} \exp\left(i \frac{x^2 + (x')^2}{2t}\right) J_\nu\left(\frac{xx'}{t}\right) \quad (t > 0, x, x' \in U) \tag{5.8}$$

(see [37, 46, 48]).

Proposition 5.1 *Let $\mathcal{T} =]0, +\infty[\times]0, +\infty[$, $\mathcal{H} : x \in]0, +\infty[\mapsto -(\partial^2/\partial x^2 - u/x^2)/2$ for some physical constant $u > 0$. For any $\lambda \in \mathbb{R}$, the initial datum $x \in]0, +\infty[\mapsto e^{i\lambda x}$ evolves through the Cauchy–Kowalevski Schrödinger equation (5.6) to a function $(t, x) \mapsto \varphi_\lambda(t, x)$ which is C^∞ in \mathcal{T} . For any $a \in \mathbb{R} \setminus [-1, 1]$, the family $\{\varphi_\lambda; \lambda \in \mathbb{R}\}$ admits as a \mathcal{F} -supershift (in the sense of Definition 5.1) the sequence*

$$\{(t, x) \in \mathcal{T} \mapsto \psi_N(t, x, a)\}_{N \geq 1} = \left\{ \sum_{j=0}^N C_j(N, a) \varphi_{1-2j/N} \right\}$$

with \mathcal{F} -supershift domain equal to \mathcal{T} . Moreover, for any $(\mu, \nu) \in \mathbb{N}^2$, the sequence of functions

$$\frac{\partial^{\mu+\nu}}{\partial t^\mu \partial x^\nu}(\psi_N(t, x, a)) = \frac{1}{(2i)^\mu} \left(\left(-\frac{\partial^2}{\partial x^2} + \frac{u}{x^2} \right)^{\odot \mu} \odot \frac{\partial^\nu}{\partial x^\nu} \right) (\psi_N(t, x, a))$$

converges uniformly on any compact $K \subset \subset \mathcal{T}$ to the function

$$(t, x) \in \mathcal{T} \mapsto \frac{1}{(2i)^\mu} \left(\left(-\frac{\partial^2}{\partial x^2} + \frac{u}{x^2} \right)^{\odot \mu} \odot \frac{\partial^\nu}{\partial x^\nu} \right) (\varphi_a(t, x)).$$

Proof Let $\lambda \in \mathbb{R}$. The evolution of the initial datum $x \in U \mapsto e^{i\lambda x}$ through the Schrödinger equation (5.6) is explicited (for the moment formally) thanks to the expression (5.8) of the Green function as

$$(t, x) \in \mathcal{T} \mapsto \frac{(-i)^{\nu+1}}{2^\nu} e^{ix^2/(2t)} \frac{x^{\nu+1/2}}{t^{\nu+1}} \int_0^\infty (x')^{\nu+1/2} e^{i(x')^2/(2t)} E_\nu\left(\frac{xx'}{t}\right) e^{i\lambda x'} dx'. \tag{5.9}$$

For any $M \in \mathbb{N}$ such that $2M > \nu - 1/2$ and any $y > 0$, one has

$$E_\nu(y) = \frac{1}{\sqrt{\pi}} \left(\frac{2}{y}\right)^{\nu+1/2} \left(\cos(y - \nu\pi/4 - \pi/2) \left(\sum_{\kappa=0}^{M-1} (-1)^\kappa \frac{a_{2\kappa}(\nu)}{y^{2\kappa}} + R_{2M}(\nu, y) \right) \right. \\ \left. + \sin(y - \nu\pi/4 - \pi/2) \left(\sum_{\kappa=0}^{M-1} (-1)^\kappa \frac{a_{2\kappa+1}(\nu)}{y^{2\kappa+1}} + R_{2M+1}(\nu, y) \right) \right)$$

with

$$|R_{2M}(\nu, y)| < \frac{|a_{2M}(\nu)|}{y^{2M}}, \quad |R_{2M+1}(\nu, y)| < \frac{|a_{2M+1}(\nu)|}{y^{2M+1}},$$

where

$$a_k(\nu) = (-1)^k \frac{\cos(\pi\nu)}{\pi} \frac{\Gamma(k + 1/2 + \nu) \Gamma(k + 1/2 - \nu)}{2^k \Gamma(k + 1)} \quad \forall k \in \mathbb{N}$$

(see [49, pp. 207–209]). It follows from such developments, together with Proposition 4.1, that the integral in (5.9) exists for any $(t, x) \in \mathcal{T}$ as a semi-convergent integral (of the Fresnel-type), whose value coincides with the regularized integral described in Sect. 4. Set now

$$\mathfrak{T} =]0, +\infty[, \phi : t \in \mathfrak{T} \mapsto -\frac{1}{2t} \in]-\infty, 0[\\ B_j : (t, Z, \check{Z}) \in \mathfrak{T} \times \mathbb{C}^2 \mapsto \begin{cases} E_\nu\left(\frac{Z\check{Z}}{t}\right) & \text{if } j = 0, \\ 0 & \text{if } j \in \mathbb{N}^*. \end{cases} \quad \chi := \nu + \frac{1}{2},$$

in order to fit with the setting described in Theorem 5.1. Since $E_\nu \in A_1(\mathbb{C})$ and

$$\frac{|Z \check{Z}|}{t} = \frac{1}{t} \times \varepsilon |Z| \times \frac{|\check{Z}|}{\varepsilon} \leq \frac{1}{2t} \left(\varepsilon^2 |Z|^2 + \frac{|\check{Z}|^2}{\varepsilon^2} \right) \forall t > 0, \forall (Z, \check{Z}) \in \mathbb{C}^2$$

the operator with order 0 given as $t \mapsto B_0(t, Z, \check{Z}) (d/dZ)^0$ satisfies the hypothesis of this theorem with $p = 1$ and $\check{p} = 2$. Then the operator

$$\mathbb{D}(t) = \int_0^\infty Z^{\nu+1/2} e^{-i Z^2/(2t)} E_\nu \left(\frac{Z \cdot \check{Z}}{t} \right) (\cdot) dZ$$

acts continuously locally uniformly in $t \in]0, +\infty[$ from $A_1(\mathbb{C})$ into $A_2(\mathbb{C})$. For any $\lambda \in \mathbb{R}$ and $t > 0$, the function $x \in]0, +\infty[\mapsto \varphi_\lambda(t, x)$ is C^∞ because of its expression (5.9). Moreover, when $a \in \mathbb{R} \setminus [-1, 1]$, it follows from Theorem 2.1 that the sequence

$$\left\{ z \in \mathbb{C} \mapsto \mathbb{D}(t) \left(\sum_{j=0}^N C_j(N, a) e^{i(1-2j/N)(\cdot)} \right) (z) \right\}_{N \geq 1}$$

converges in $A_2(\mathbb{C})$ (locally uniformly with respect to $t > 0$) to $z \mapsto \mathbb{D}(t)(e^{ia(\cdot)})$. One concludes then to the second assertion in the statement of the theorem. As for the last assertion, it follows from the fact that the action of $i\partial/\partial t$ and $\mathcal{H}(x)$ coincide on solutions of (5.6), together with the continuity of the operator d/dz from $A_2(\mathbb{C})$ into itself. □

5.4 The Schrödinger Cauchy Problem for the Quantum Harmonic Oscillator

Let now $U = \mathbb{R}$ and the hamiltonian in (5.6) be $x \in \mathbb{R} \mapsto \mathcal{H}(x) = -(\partial^2/\partial x^2)/2 + x^2/2$. The corresponding Cauchy–Kowalevski problem (with $]0, +\infty[\times \mathbb{R}$ as phase space) is the *Schrödinger Cauchy problem for the quantum harmonic oscillator*, see [13, §5.3, §. 6.4] or [31] for more references. In this case again, the Green function can be explicitated and is therefore handable. It is the locally integrable function in $]0, +\infty[\times \mathbb{R} \times \mathbb{R}$ defined as

$$\begin{aligned} G(t, x, t' = 0, x') &= \sqrt{\frac{1}{2i\pi \sin t}} e^{i \left(\frac{(x^2 + (x')^2) \cos t - 2xx'}{2 \sin t} \right)} \\ &= \left(\sqrt{\frac{1}{2i\pi \sin t}} e^{i \frac{\cotan t}{2} x^2} \right) e^{i \frac{\cotan t}{2} (x')^2} e^{-i \frac{xx'}{\sin t}} \quad (t > 0, x, x' \in \mathbb{R}). \end{aligned} \tag{5.10}$$

Proposition 5.2 *Let $\mathcal{T} =]0, +\infty[\times \mathbb{R}$ and $\mathcal{H} : x \in \mathbb{R} \mapsto -(\partial^2/\partial x^2 - x^2)/2$. For any $\lambda \in \mathbb{R}$, the initial datum $x \in \mathbb{R} \mapsto e^{i\lambda x}$ evolves through the Cauchy–Kowalevski Schrödinger equation (5.6) to a \mathbb{C} -valued distribution $\varphi_\lambda \in \mathcal{D}'(\mathcal{T}, \mathbb{C})$ with singular support $\pi(2\mathbb{N}+1)/2 \times \mathbb{R}$. Let $\mathcal{U} = \mathcal{T} \setminus (\pi(2\mathbb{N}+1)/2 \times \mathbb{R})$. For any $a \in \mathbb{R} \setminus [-1, 1]$, the*

family $\{(\varphi_\lambda)|_{\mathcal{U}}; \lambda \in \mathbb{R}\}$, considered as a family of functions, admits as a \mathcal{F} -supershift (in the sense of Definition 5.1) the sequence

$$\{(t, x) \in \mathcal{U} \mapsto \psi_N(t, x, a)\}_{N \geq 1} = \left\{ \sum_{j=0}^N C_j(N, a) (\varphi_{1-2j/N})|_{\mathcal{U}} \right\}_{N \geq 1}$$

with \mathcal{F} -supershift domain equal to \mathcal{U} . Moreover, for any $(\mu, \nu) \in \mathbb{N}^2$, the sequence of functions from \mathcal{U} to \mathbb{C}

$$\frac{\partial^{\mu+\nu}}{\partial t^\mu \partial x^\nu}(\psi_N(t, x, a)) = \frac{1}{(2i)^\mu} \left(\left(-\frac{\partial^2}{\partial x^2} + x^2 \right)^{\odot \mu} \odot \frac{\partial^\nu}{\partial x^\nu} \right) (\psi_N(t, x, a))$$

converges uniformly on any compact $K \subset \subset \mathcal{U}$ to the function

$$(t, x) \in \mathcal{F} \mapsto \frac{1}{(2i)^\mu} \left(\left(-\frac{\partial^2}{\partial x^2} + x^2 \right)^{\odot \mu} \odot \frac{\partial^\nu}{\partial x^\nu} \right) (\varphi_a(t, x)).$$

Proof Consider the two (for the moment formal) operators

$$t \in]0, +\infty[\setminus \pi\mathbb{N}^*/2 \mapsto |\sin t| \int_0^\infty e^{i \frac{\sin 2t}{4} Z^2} e^{-i\varpi \operatorname{sign}(\sin(t)) \check{Z} Z} \circ_{H_{|\sin t|}(\cdot)} dZ \tag{5.11}$$

($\varpi = \pm 1$) which appear (after performing the change of variables $Z \leftrightarrow |\sin t| Z$ on $]0, +\infty[$) in the splitting of

$$t \in]0, +\infty[\setminus \pi\mathbb{N}^*/2 \mapsto \int_{\mathbb{R}} e^{i \frac{\cotan t}{2} Z^2} e^{-i \check{Z} Z / \sin t} (\cdot) dZ$$

(see Remark 4.1). Set now

$$\begin{aligned} \mathfrak{T} &=]0, +\infty[\setminus \pi\mathbb{N}^*/2, \quad \phi : t \in \mathfrak{T} \mapsto -\frac{\sin(2t)}{4} \\ B_j : (t, Z, \check{Z}) \in \mathfrak{T} \times \mathbb{C}^2 &\mapsto \begin{cases} \exp(-i \varpi \operatorname{sign}(\sin(t)) Z \check{Z}) \odot H_{|\sin t|} & \text{if } j = 0, \\ 0 & \text{if } j \in \mathbb{N}^*. \end{cases} \quad \chi = 0 \end{aligned}$$

($\varpi = \pm 1$) in order to fit with the setting described in Theorem 5.1. As in the proof of Proposition 5.1, this theorem applies here and the two operators (5.11) act continuously from $A_1(\mathbb{C})$ to $A_2(\mathbb{C})$ (locally uniformly with respect to the parameter $t \in \mathfrak{T}$). Note again that the Fresnel-type integrals (5.11), where $Z \mapsto e^{i\lambda Z}$ ($\lambda \in \mathbb{R}$) is taken inside the bracket and $\check{Z} \in \mathbb{R}$, are semi-convergent and their values as semi-convergent integrals coincide with the values that are obtained by regularization as in

Sect. 4. In fact, in the case where $\check{Z} = x \in \mathbb{R}$ and $t \in \mathfrak{T}$, the value of

$$\left(\sqrt{\frac{1}{2i\pi \sin t}} e^{i \frac{\cotan t}{2} \check{Z}^2} \right) \int_{\mathbb{R}} e^{i \frac{\cotan t}{2} Z^2} e^{-i \check{Z}Z / \sin t} (\cdot) dZ$$

(understood as a regularized integral, see Sect. 4, in particular Remark 4.1) equals

$$(\cos t)^{-1/2} e^{-i \check{Z}^2 \frac{\tan(t)}{2}} e^{-i \lambda^2 \frac{\tan(t)}{2}} \odot H_{1/\cos t}(e^{i\lambda(\cdot)})(\check{Z})$$

(see [13, Proposition 5.3.1]). Since $(t, x) \in]0, +\infty[\times \mathbb{R} \mapsto (\cos t)^{-1/2} e^{-ix^2 \tan(t)} e^{-i\lambda^2 \tan(t)}$ is a locally integrable function, the initial datum $x \in \mathbb{R} \mapsto e^{i\lambda x}$ evolves through the Schrödinger equation (5.6) as a distribution φ_λ (in fact defined by a locally integrable function). Let $\mathbb{D}(t)$ the differential operator

$$\mathbb{D} : t \in]0, \infty[\setminus \pi \frac{2\mathbb{N} + 1}{2} \mapsto \sum_{j=0}^{\infty} \frac{1}{j!} \left(i \frac{\sin 2t}{4} \right)^j (d/dW)^{2j}.$$

Since

$$\begin{aligned} & (\cos t)^{-1/2} e^{-i \check{Z}^2 \frac{\tan(t)}{2}} e^{-i \lambda^2 \frac{\tan(t)}{2}} \odot H_{1/\cos t}(e^{i\lambda(\cdot)})(\check{Z}) \\ &= (\cos t)^{-1/2} e^{-i \check{Z}^2 \frac{\tan(t)}{2}} \mathbb{D}(t)(e^{i\lambda(\cdot)})(\check{Z}), \end{aligned}$$

and \mathbb{D} acts continuously locally uniformly in t from $A_1(\mathbb{C})$ to $A_2(\mathbb{C})$ thanks to Lemma 2.2, the sequence

$$\left\{ \sum_{j=0}^N C_j(N, a) (\varphi_{1-2j/N})|_{\mathcal{U}} \right\}_{N \geq 1}$$

is, for any $a \in \mathbb{R} \setminus [-1, 1]$, a supershift for the family $\mathcal{F} = \{(\varphi_\lambda)|_{\mathcal{U}}; \lambda \in \mathbb{R}\}$ (with \mathcal{F} -supershift domain \mathcal{U}). The last assertion follows from the same argument than that used for the last assertion in Proposition 5.1. □

6 Singularities in the Quantum Harmonic Oscillator Evolution

This section is the natural continuation of Sect. 5.4. We continue to investigate with respect to the notion of supershift the evolution of initial data $x \in \mathbb{R} \mapsto e^{i\lambda x}$, when $\lambda \in \mathbb{R}$, through the Cauchy–Schrödinger problem for the quantum harmonic oscillator and focus now on singularities. In this section we keep the same notations as in Proposition 5.2 and fix a point (t_0, x_0) in $\mathcal{T} \setminus \mathcal{U}$. We will just consider the case $t_0 = \pi/2$ since the situation is essentially identical at any point $((2k + 1)\pi/2, x_0)$ with $k \in \mathbb{N}$ and $x_0 \in \mathbb{R}$.

Let, for $\lambda \in \mathbb{R}$, $\varphi_\lambda \in \mathcal{D}'(\mathcal{T}, \mathbb{C})$ be the distribution evolved from the initial datum $x \mapsto e^{i\lambda x}$ through the Schrödinger operator for the quantum oscillator problem (5.6) (with $\mathcal{H} : x \in \mathbb{R} \mapsto (-\partial^2/\partial x^2 + x^2)/2$).

Let $\theta \in \mathcal{D}(\mathcal{T}, \mathbb{C})$ be a test-function with support in a small neighborhood of $(\pi/2, x_0)$ and $(t, x) \mapsto \xi(t, x) := \theta(t, x) \exp((ix^2 \cotan t)/2)/\sqrt{2i\pi}$. One has (formally) for any $\lambda \in \mathbb{R}$,

$$\begin{aligned} \langle \varphi_\lambda, \theta \rangle &= - \int_{\mathbb{R}^2} \left[\int_{\mathbb{R}} e^{i \frac{\sin u}{2} Z^2} e^{-i \check{Z} Z / \sqrt{\cos u}} (e^{i\lambda(\cdot)}) dZ \right]_{\check{Z}=x} \xi(\pi/2 - u, x) du dx \\ &= \int_{\mathbb{R}^2} \left[\int_{\mathbb{R}} e^{i \frac{\sin u}{2} Z^2} e^{-i \check{Z} Z} (e^{i\lambda(\cdot)}) dZ \right]_{\check{Z}=x} \tilde{\xi}(u, x) du dx, \end{aligned} \tag{6.1}$$

where $\tilde{\xi}(u, x) = -\sqrt{\cos u} \xi(\pi/2 - u, \sqrt{\cos u} x)$ is a test-function with support about $(0, x_0)$. The regularized integral is then

$$\begin{aligned} &\lim_{\varepsilon \rightarrow 0+} \int_{\mathbb{R}^2} \left[\int_{\mathbb{R}} e^{-\varepsilon Z^2} e^{i \frac{\sin u}{2} Z^2} e^{-i \check{Z} Z} (e^{i\lambda(\cdot)}) dZ \right]_{\check{Z}=x} \tilde{\xi}(u, x) du dx \\ &= \lim_{\varepsilon \rightarrow 0+} \int_{\mathbb{R}^2} \left[\int_{\mathbb{R}} e^{-\varepsilon Z^2} e^{i \frac{\sin u}{2} (Z^2 - 2(\check{Z} - \lambda) / \sin u)} dZ \right]_{\check{Z}=x} \tilde{\xi}(u, x) du dx \\ &= \int_{\mathbb{R}^2} \left[\exp\left(\frac{2i}{\sin u} (\check{Z} - \lambda)^2\right) \right]_{\check{Z}=x} \sqrt{\frac{2i\pi}{\sin u}} \tilde{\xi}(u, x) du dx \\ &= \int_{\mathbb{R}^2} \left[\exp\left(\frac{i}{v} (\check{Z} - \lambda)^2\right) \right]_{\check{Z}=x} \sqrt{\frac{1}{v}} \tilde{\theta}(v, x) dv dx \end{aligned}$$

for some test-function $(v, x) \mapsto \tilde{\theta}(v, x)$ with support about $(0, x_0)$ (one uses here Lebesgue domination theorem and the change of variables $(\sin u)/2 \longleftrightarrow v$ about $u = 0$). Though such expression makes sense when $\lambda \in \mathbb{R}$ (since $|\exp(i(x - \lambda)^2/v)| = 1$ for any point $(v, x) \in \text{Supp}(\tilde{\xi})$), it does not make sense anymore when $\lambda \in \mathbb{C}$. In order to overcome this difficulty, one needs to formulate the following lemma.

Lemma 6.1 *Let $\mathbb{D}(\check{Z})$ ($\check{Z} \in \mathbb{C}$) be a differential operator of the form*

$$\sum_{\kappa=0}^{\infty} \left[\frac{A_\kappa(\check{Z}, (d/dZ))}{\kappa!} (\cdot) \right]_{Z=0} (d/dv)^\kappa, \tag{6.2}$$

(where $A_\kappa \in \mathbb{C}[[\check{Z}, d/dZ]]$ for any $\kappa \in \mathbb{N}$), considered as acting from the space of entire functions of the variable Z to the space $\mathbb{C}[[\check{Z}]][[d/dv]]$. Suppose that there exist $p \geq 1$ and $\check{p} \geq 1$ and $B, \check{B} \geq 0$ such that

$$\sup_{\kappa \in \mathbb{N}, \check{Z} \in \mathbb{C}} (|A_\kappa(\check{Z}, W)| \exp(-B|W|^p - \check{B}|W|^{\check{p}})) < +\infty. \tag{6.3}$$

Then, for any $b \geq 0$, there exists $A^{(b)} \geq 0$ such that

$$\forall C \geq 0, \quad \forall f \in A_1^{C,b}(\mathbb{C}), \quad \sup_{\kappa \in \mathbb{N}} |A_\kappa(\check{Z}, (d/dZ))(f)(0)| \leq C A^{(b)} e^{\check{B}|\check{Z}|^{\check{p}}}.$$

In particular, for any $f \in A_1^{C,b}(\mathbb{C})$, $\mathbb{D}(\check{Z})(f)$ remains an infinite order differential operator $\sum_{\kappa \geq 0} \alpha_\kappa(\check{Z})(f) (d/dv)^\kappa$ with coefficients satisfying (independently of $f \in A_1^{C,b}(\mathbb{C})$)

$$\sum_{\kappa \in \mathbb{N}} k! |\alpha_\kappa(\check{Z})(f)| \exp(-B|\check{Z}|^{\check{p}}) = C A^{(b)} < +\infty.$$

Proof The coefficients of A_κ as a polynomial in d/dZ satisfy

$$\sum_{\kappa, j \in \mathbb{N}, \check{Z} \in \mathbb{C}} |a_{\kappa, j}(\check{Z})| \leq C_0 \frac{b_0^j}{\Gamma(j/p) + 1} e^{\check{B}|\check{Z}|^{\check{p}}}$$

for some absolute constants C_0 and b_0 (Lemma 2.1). As in the proof of Lemma 2.2, one concludes that for any $f \in A_1^{C,b}(\mathbb{C})$ and any $\kappa \in \mathbb{N}$, one has uniform estimates $|A_\kappa(\check{Z}, d/dZ)(f)| \leq C A^{(b)} \exp(b_0 b |Z| + \check{B}|\check{Z}|^{\check{p}})$ for some positive constant $A^{(b)}$. One gets the required estimates when evaluating at $Z = 0$. \square

One can then complete Proposition 5.2 into the following companion proposition.

Proposition 6.1 *Let $\mathcal{T} =]0, +\infty[\times \mathbb{R}$ and $\mathcal{H} : x \in \mathbb{R} \mapsto -(\partial^2/\partial x^2 + x^2)/2$. For any $\lambda \in \mathbb{R}$, let $\varphi_\lambda \in \mathcal{D}'(\mathcal{T}, \mathbb{C})$ be the evolved distribution from the initial datum $x \in \mathbb{R} \mapsto e^{i\lambda x}$ through the Cauchy–Kowalevski Schrödinger equation (5.6). Let $\mathcal{F} = \{\varphi_\lambda; \lambda \in \mathbb{R}\}$, where each φ_λ is considered as a hyperfunction in \mathcal{T} . Then, for any $a \in \mathbb{R} \setminus [-1, 1]$, the sequence $\{\sum_{j=0}^N C_j(N) \varphi_{1-2j/N}\}_{N \geq 1}$ is a \mathcal{F} -supershift of hyperfunctions over the \mathcal{F} -supershift domain \mathcal{T} .*

Proof Let $\theta \in \mathcal{D}(\mathbb{R}_{t,x}^2, \mathbb{C})$ with support a small neighborhood V of the point $(\pi/2, x_0)$ ($x_0 \in \mathbb{R}$) and $\tilde{\theta}$ the test-function with support $V - (\pi/2, 0) \ni (0, x_0)$ that corresponds to it through the successive transformations explicited previously. One has for any $\lambda \in \mathbb{R}$,

$$\begin{aligned} \langle \varphi_\lambda, \theta \rangle &= \int_{\mathbb{R}} \int_0^\infty \left[\left(\exp\left(\frac{i}{v}(\check{Z} - \lambda)^2\right) \right) \right]_{\check{Z}=x} \frac{\tilde{\theta}(v, x)}{\sqrt{v}} dv dx \\ &\quad - i \int_{\mathbb{R}} \int_0^\infty \left[\left(\exp\left(-\frac{i}{v}(\check{Z} - \lambda)^2\right) \right) \right]_{\check{Z}=x} \frac{\tilde{\theta}(-v, x)}{\sqrt{v}} dv dx \\ &= \int_{\mathbb{R}} \int_0^\infty \left(\left[\sum_{\kappa=0}^\infty \frac{i^\kappa}{\kappa!} \frac{(\check{Z} - \lambda)^{2\kappa}}{v^{1/2+\kappa}} \right]_{\check{Z}=x} \tilde{\theta}(v, x) \right. \\ &\quad \left. - i \left[\sum_{\kappa=0}^\infty \frac{(-i)^\kappa}{\kappa!} \frac{(\check{Z} - \lambda)^{2\kappa}}{v^{1/2+\kappa}} \right]_{\check{Z}=x} \tilde{\theta}(-v, x) \right) dv dx. \end{aligned} \tag{6.4}$$

For any $\kappa \in \mathbb{N}$, the distribution $v_+^{-1/2-\kappa} \in \mathcal{D}'([0, +\infty[, \mathbb{R})$ can be expressed as

$$v_+^{-1/2-\kappa} = \frac{2^\kappa}{\prod_{\ell=1}^\kappa (2(\kappa - \ell) + 1)} (-d/dv)^\kappa (v_+^{-1/2})$$

in the sense of distributions in $\mathcal{D}'([0, +\infty[, \mathbb{R})$. Then, one can reformulate formally (6.2) as

$$\begin{aligned} \langle \varphi_\lambda, \theta \rangle &= \sum_{\kappa=0}^\infty \frac{(2i)^\kappa}{\kappa! \prod_{\ell=1}^\kappa (2(\kappa - \ell) + 1)} \\ &\int_{\mathbb{R}} \left[\left(\check{Z} + i \frac{d}{dZ} \right)^{2\kappa} (e^{i\lambda(\cdot)}) \right]_{\check{Z}=x} (0) \left(\frac{d}{dv} \right)^\kappa (v_+^{-1/2}), \tilde{\theta}(\cdot, x) - i(-1)^\kappa \tilde{\theta}(-\cdot, x) \Big] dx. \end{aligned} \tag{6.5}$$

Lemma 6.1 applies to the two operators

$$\begin{aligned} \mathbb{D}(\check{Z}) &= \sum_{\kappa=0}^\infty \frac{1}{\kappa!} \left[\frac{(2i)^\kappa (\check{Z} + id/dZ)^{2\kappa}}{\prod_{\ell=1}^\kappa (2(\kappa - \ell) + 1)} (\cdot) \right]_{Z=0} (d/dv)^\kappa \\ \tilde{\mathbb{D}}(\check{Z}) &= \sum_{\kappa=0}^\infty \frac{1}{\kappa!} \left[\frac{(-i)^{\kappa+1} 2^\kappa (\check{Z} + id/dZ)^{2\kappa}}{2 \prod_{\ell=1}^\kappa ((\kappa - \ell) + 1)} (\cdot) \right]_{Z=0} (d/dv)^\kappa \end{aligned} \tag{6.6}$$

with $p = \check{p} = 2$. These two operators act then continuously (locally uniformly with respect to the parameter \check{Z}) from $A_1(\mathbb{C})$ into the space of infinite order differential operators in d/dv (depending on the parameter $\check{Z} \in \mathbb{C}$). Such differential operators can be considered as hyperfunctions on \mathbb{R}_v (elements of $\mathcal{H}(\mathbb{R}_v)$). Since $v_+^{-1/2}$ is a Fourier hyperfunction in the real line \mathbb{R} , the two $\mathcal{H}(\mathbb{R})$ -valued operators $f \in A_1(\mathbb{C}) \mapsto \mathbb{D}(\check{Z})(f) \odot v_+^{-1/2}$ and $f \in A_1(\mathbb{C}) \mapsto \tilde{\mathbb{D}}(\check{Z})(f) \odot v_+^{-1/2}$ are well defined (see [40, Proposition 8.4.8 and Exercise 8.4.5]) and depend continuously (locally uniformly with respect to \check{Z}) on the entry f in $A_1(\mathbb{C})$. Proposition 6.1 follows then from Theorem 2.1 and from the expression (6.4) (together with its formal reformulation (6.5)) for the evaluations $\langle \varphi_\lambda, \theta \rangle$ when $\lambda \in \mathbb{R}$ and φ_λ is considered as an element in $\mathcal{D}'(\mathcal{T}, \mathbb{C})$ (acting on $\theta \in \mathcal{D}(\mathcal{T}, \mathbb{C})$) which can be also interpreted as a hyperfunction on \mathcal{T} . □

Data availability statement There are no data used in this manuscript.

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References

1. Aharonov, Y., Behrndt, J., Colombo, F., Schlosser, P.: Schrödinger evolution of superoscillations with δ - and δ' -potentials. *Quantum Stud. Math. Found.* **7**(3), 293–305 (2020)
2. Aharonov, Y., Behrndt, J., Colombo, F., Schlosser, P.: Green's function for the Schrödinger equation with a generalized point interaction and stability of superoscillations. *J. Differ. Equ.* **277**, 153–190 (2021)
3. Aharonov, Y., Behrndt, J., Colombo, F., Schlosser, P.: A unified approach to Schrödinger evolution of superoscillations and supershifts. [arXiv:2102.11795](https://arxiv.org/abs/2102.11795)
4. Aharonov, Y., Colombo, F., Struppa, D.C., Tollaksen, J.: Schrödinger evolution of superoscillations under different potentials. *Quantum Stud. Math. Found.* **5**, 485–504 (2018)
5. Aharonov, Y., Colombo, F., Sabadini, I., Struppa, D.C., Tollaksen, J.: Evolution of superoscillations in the Klein–Gordon field. *Milan J. Math.* **88**(1), 171–189 (2020)
6. Aharonov, Y., Sabadini, I., Tollaksen, J., Yger, A.: Classes of superoscillating functions. *Quantum Stud. Math. Found.* **5**, 439–454 (2018)
7. Aharonov, Y., Albert, D., Vaidman, L.: How the result of a measurement of a component of the spin of a spin-1/2 particle can turn out to be 100. *Phys. Rev. Lett.* **60**, 1351–1354 (1988)
8. Aharonov, Y., Colombo, F., Nussinov, S., Sabadini, I., Struppa, D.C., Tollaksen, J.: Superoscillation phenomena in $SO(3)$. *Proc. R. Soc. A* **468**, 3587–3600 (2012)
9. Aharonov, Y., Colombo, F., Sabadini, I., Struppa, D.C., Tollaksen, J.: On the Cauchy problem for the Schrödinger equation with superoscillatory initial data. *J. Math. Pures Appl.* **99**, 165–173 (2013)
10. Aharonov, Y., Colombo, F., Sabadini, I., Struppa, D.C., Tollaksen, J.: Some mathematical properties of superoscillations. *J. Phys. A* **44**, 365304 (2011)
11. Aharonov, Y., Colombo, F., Sabadini, I., Struppa, D.C., Tollaksen, J.: Superoscillating sequences as solutions of generalized Schrödinger equations. *J. Math. Pures Appl.* **103**, 522–534 (2015)
12. Aharonov, Y., Colombo, F., Sabadini, I., Struppa, D.C., Tollaksen, J.: Superoscillating sequences in several variables. *J. Fourier Anal. Appl.* **22**, 751–767 (2016)
13. Aharonov, Y., Colombo, F., Sabadini, I., Struppa, D.C., Tollaksen, J.: The Mathematics of Superoscillations. *Mem. Amer. Math. Soc.*, vol. 247, no. 1174, v+107 pp (2017)
14. Aharonov, Y., Rohrlich, D.: *Quantum Paradoxes: Quantum Theory for the Perplexed*. Wiley, Weinheim (2005)
15. Aharonov, Y., Vaidman, L.: Properties of a quantum system during the time interval between two measurements. *Phys. Rev. A* **41**, 11–20 (1990)
16. Alpay, D., Colombo, F., Sabadini, I., Struppa, D.C.: Aharonov–Berry superoscillations in the radial harmonic oscillator potential. *Quantum Stud. Math. Found.* **7**, 269–283 (2020)
17. Alpay, D., Colombo, F., Sabadini, I.: Superoscillations and analytic extension in Schur analysis. *J. Fourier Anal. Appl.* **27**(2), 1–9 (2021)
18. Aoki, T., Colombo, F., Sabadini, I., Struppa, D.C.: Continuity theorems for a class of convolution operators and applications to superoscillations. *Ann. Mat. Pura Appl.* **197**, 1533–1545 (2018)
19. Aoki, T., Colombo, F., Sabadini, I., Struppa, D.C.: Continuity of some operators arising in the theory of superoscillations. *Quantum Stud. Math. Found.* **5**, 463–476 (2018)
20. Behrndt, J., Colombo, F., Schlosser, P.: Evolution of Aharonov–Berry superoscillations in Dirac δ -potential. *Quantum Stud. Math. Found.* **6**, 279–293 (2019)
21. Berenstein, C.A., Gay, R.: *Complex Variables. An Introduction*, Graduate Texts in Mathematics, vol. 125. Springer, New York (1991)
22. Berenstein, C.A., Gay, R.: *Complex Analysis and Special Topics in Harmonic Analysis*. Springer, New York (1995)
23. Berry, M.V.: Evanescent and real waves in quantum billiards and Gaussian beams. *J. Phys. A* **27**, 391 (1994)
24. Berry, M.: Exact nonparaxial transmission of subwavelength detail using superoscillations. *J. Phys. A* **46**, 205203 (2013)
25. Berry, M.V.: Faster than Fourier, in quantum coherence and reality. In: *Celebration of the 60th Birthday of Yakir, Aharonov*, pp. 55–65. J.S. Anandan and J. L. Safko, World Scientific, Singapore (1994)
26. Berry, M.V.: Representing superoscillations and narrow Gaussians with elementary functions. *Milan J. Math.* **84**, 217–230 (2016)
27. Berry, M., Dennis, M.R.: Natural superoscillations in monochromatic waves in D dimension. *J. Phys. A* **42**, 022003 (2009)

28. Berry, M.V., Popescu, S.: Evolution of quantum superscillations, and optical superresolution without evanescent waves. *J. Phys. A* **39**, 6965–6977 (2006)
29. Berry, M.V., Shukla, P.: Pointer supershifts and superscillations in weak measurements. *J. Phys. A* **45**, 015301 (2012)
30. Berry, M.V., et al.: Roadmap on superscillations. *J. Opt.* **21**(5), 053002 (2019)
31. Buniy, R., Colombo, F., Sabadini, I., Struppa, D.C.: Quantum harmonic oscillator with superscillating initial datum. *J. Math. Phys.* **55**, 113511 (2014)
32. Colombo, F., Gantner, J., Struppa, D.C.: Evolution by Schrödinger equation of Aharonov–Berry superscillations in centrifugal potential. *Proc. R. Soc. A* **475**(2225), 20180390 (2019)
33. Colombo, F., Struppa, D.C., Yger, A.: Superscillating sequences towards approximation in \mathcal{S}' -type spaces and extrapolation. *J. Fourier Anal. Appl.* **25**, 242–266 (2019)
34. Colombo, F., Sabadini, I., Struppa, D.C., Yger, A.: Superscillating sequences and hyperfunctions. *Publ. Res. Inst. Math. Sci.* **55**, 665–688 (2019)
35. Cordero-Soto, R., Suazo, E., Suslov, S.K.: Quantum integrals of motion for variable quadratic Hamiltonians. *Ann. Phys.* **325**(9), 1884–1912 (2010)
36. Ferreira, P.J.S.G., Kempf, A.: Superscillations: faster than the Nyquist rate. *IEEE Trans. Signal Process.* **54**, 3732–3740 (2006)
37. Ferreira, P.J.S.G., Kempf, A.: Unusual properties of superscillating particles. *J. Phys. A* **37**, 12067–76 (2004)
38. Gorenflo, R., Kilbas, A.A., Mainardi, F., Rogosin, S.V.: Mittag–Leffler Functions, Related Topics and Applications. Springer Monographs in Mathematics, Springer, Heidelberg (2014)
39. Gradshteyn, I.S., Ryzhik, I.M.: Table of Integrals, Series, and Products, 7th edn. Elsevier/Academic Press, Amsterdam (2007)
40. Kaneko, A.: Introduction to Hyperfunctions. Kluwer, Mathematics and Its Applications (1988)
41. Kempf, A.: Four aspects of superscillations. *Quantum Stud. Math. Found.* **5**, 477–484 (2018)
42. Khandekar, K.C., Lawande, S.V.: Exact propagator for a time-dependent harmonic oscillator with and without a singular perturbation. *J. Math. Phys.* **16**, 384 (1975)
43. Lee, D.G., Ferreira, P.J.S.G.: Superscillations of prescribed amplitude and derivative. *IEEE Trans. Signal Process.* **62**, 3371–3378 (2014)
44. Lee, D.G., Ferreira, P.J.S.G.: Superscillations with optimal numerical stability. *IEEE Signal Process. Lett.* **21**(12), 1443–1447 (2014)
45. Lindberg, J.: Mathematical concepts of optical superresolution. *J. Opt.* **14**, 083001 (2012)
46. Schulman, L.S.: Techniques and Applications of Path Integration. Wiley, New York (1981)
47. Taylor, B.A.: Some locally convex spaces of entire functions. In: Korevaar, J., Chern, S.S., Ehrenpreis, L., Fuchs, W.H.J., Rubel, L.A. (eds.) Entire Functions and Related Parts of Analysis, Proceedings of Symposia in Pure Mathematics, vol. 11, pp. 431–467. American Mathematical Society (1968)
48. Tsauro, G.Y., Wang, J.: Constructing Green functions of the Schrödinger equation by elementary transformations. *Am. J. Phys.* **74**(7), 600–606 (2006)
49. Watson, G.N.: A Treatise on the Theory of Bessel Functions, Reprint of the second (1944) edition. Cambridge Mathematical Library. Cambridge University Press, Cambridge (1995)